



Equivalence of weak and viscosity solutions for the nonhomogeneous double phase equation

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Abstract

We establish the equivalence between weak and viscosity solutions to the nonhomogeneous double phase equation with lower-order term

$$-\operatorname{div}(|Du|^{p-2}Du + a(x)|Du|^{q-2}Du) = f(x, u, Du), \quad 1 < p \leq q < \infty, \quad a(x) \geq 0.$$

We find some appropriate hypotheses on the coefficient $a(x)$, the exponents p, q and the nonlinear term f to show that the viscosity solutions with *a priori* Lipschitz continuity are weak solutions of such equation by virtue of the inf(sup)-convolution techniques. The reverse implication can be concluded through comparison principles. Moreover, we verify that the bounded viscosity solutions are exactly Lipschitz continuous, which is also of independent interest.

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1 Introduction

Let Ω be a bounded domain in \mathbb{R}^n ($n \geq 2$). In this work we aim to examine the inner relationship between weak and viscosity solutions to the following nonhomogeneous double phase equation

$$-\operatorname{div}(|Du|^{p-2}Du + a(x)|Du|^{q-2}Du) = f(x, u, Du) \quad \text{in } \Omega, \quad (1.1)$$

where $1 < p \leq q < \infty$, $a(x) \geq 0$ and $f(x, \tau, \xi) : \Omega \times \mathbb{R} \times \mathbb{R}^n \rightarrow \mathbb{R}$ is a continuous function. The double phase problems, stemming from the models of strongly anisotropic materials, were originally investigated by Zhikov [51, 52] and Marcellini [40] in the context of homogenization and Lavrentiev phenomenon.

Over the last years, problems of the type considered in (1.1) have attracted intensive attention from the variational point of view, whose celebrated prototype is given by the following unbalanced energy functional

$$W^{1,1}(\Omega) \ni u \mapsto \mathcal{P}(u, \Omega) := \int_{\Omega} (|Du|^p + a(x)|Du|^q) dx.$$

The significant characteristics of this functional are that its ellipticity and growth rate will change drastically according to the modulating coefficient $a(\cdot)$ equal to 0 or not. The regularity of minimizers is determined via a delicate interaction between the growth conditions and the pointwise behaviour of $a(\cdot)$. For instance, under the hypotheses that

$$0 \leq a(\cdot) \in C^{0,\alpha}(\Omega), \alpha \in (0, 1] \quad \text{and} \quad \frac{q}{p} \leq 1 + \frac{\alpha}{n}, \quad (1.2)$$

Colombo, Mingione *et. al.* [6, 7, 13] established the gradient Hölder continuity and Harnack inequality for the minimizers of \mathcal{P} . A key feature of this problem is that the minimizers could be even discontinuous when condition (1.2) is violated, by means of the counterexamples presented in [21, 27]. For the double phase equation

$$-\operatorname{div}(|Du|^{p-2}Du + a(x)|Du|^{q-2}Du) = -\operatorname{div}(|F|^{p-2}F + a(x)|F|^{q-2}F) \quad \text{in } \Omega,$$

the Calderón-Zygmund estimates of weak solutions were derived in [14, 18] under assumption (1.2) (see also [4, 9]). More recently, De Filippis and Mingione [19] considered a very large class of vector-valued nonautonomous variational problems involving integral functionals of the double phase type, where the authors provided a comprehensive treatment of Lipschitz regularity of solutions under sharp conditions. Despite their relatively short history, double phase problems have achieved very fruitful results with several connections to other aspects, such as the existence and multiplicity of solutions [47], the nonlocal version [20, 22], the properties of eigenvalues and

eigenfunctions [12, 46], as well as the removability and obstacle problems [10, 35]. We also refer to [5, 17, 25, 26, 29, 37, 41, 44] and references therein for more results.

The topic on equivalence of different solutions started from the works of Lions [39] and Ishii [30] on linear equations. For what concerns the quasilinear case, Juutinen, Lindqvist and Manfredi [33] proved that the weak solutions coincide with the viscosity solutions to the p -Laplace equation and its parabolic version based on the uniqueness machinery of solutions; see [34] for $p(x)$ -Laplace type equation. The equivalence of solutions was generalized to the fractional p -Laplace equation in [36] by following analogous ideas. Julin and Juutinen [32] gave a more immediate proof for the equivalence of viscosity and weak solutions to the p -Laplace equation without relying on the comparison principle of viscosity solutions. They introduced a technical regularization process through infimal convolution, which was applied to various equations incorporating the normalized $p(x)$ -Laplace equation [49], the nonhomogeneous nonlocal p -Laplace equation [8] and the normalized p -Poisson equation [3]. More related results can be found in [42, 43, 48, 50].

From the results mentioned above, we can see that the research achievements for the double phase problems mainly focus on the weak solutions from the variational perspective and there are few results concerning the relationship between viscosity and weak solutions for the general nonuniformly elliptic equations. In particular, De Filippis and Palatucci [20] showed that the bounded viscosity solutions of the nonlocal counterpart to (1.1) are locally Hölder continuous. For the homogeneous case of (1.1), Fang and Zhang [23] established the equivalence between weak and viscosity solutions by introducing $\mathcal{A}_{H(\cdot)}$ -harmonic functions that serve as a bridge. Motivated by the previous works [20, 23], our intention in the present paper is to prove the equivalence of weak and viscosity solutions for the nonhomogeneous problem (1.1). Due to the presence of the lower-order term, we cannot introduce $\mathcal{A}_{H(\cdot)}$ -harmonic functions any more and it is hard to use the full uniqueness machinery of viscosity solutions. To this end, we revisit the inf(sup)-convolution approximation, developed in [32], to verify directly that weak solutions are equivalent to viscosity solutions under some proper preconditions.

We are now in a position to state the main contributions of this manuscript. The first one establishes the following qualitative property.

Theorem 1.1 *Let $0 < a(x) \in C^1(\Omega)$ and $\frac{q}{p} \leq 1 + \frac{1}{n}$ be in force. Suppose that $f(x, \tau, \xi)$ is uniformly continuous in $\Omega \times \mathbb{R} \times \mathbb{R}^n$, decreasing in τ , Lipschitz continuous with respect to ξ and fulfilling the following growth condition*

$$|f(x, \tau, \xi)| \leq \gamma(|\tau|)(|\xi|^{p-1} + a(x)|\xi|^{q-1}) + \Phi(x), \quad (1.3)$$

where $\gamma(\cdot) \geq 0$ is continuous and $\Phi \in L_{\text{loc}}^\infty(\Omega)$. Let u be a viscosity supersolution with local Lipschitz continuity to problem (1.1) in Ω . Then u is a weak supersolution as well.

We would like to mention that the double phase operator, compared to the usual p -Laplace operator, lacks translation invariance property and exhibits two diverse growth terms owing to the presence of $a(x)$. It will lead to an additional error term $E(\varepsilon)$ in

the key Lemma 3.1 and demand an *a priori* assumption that the viscosity solution u is locally Lipschitz continuous.

The second result is about the local Lipschitz continuity of viscosity solutions, which is also of independent interest.

Theorem 1.2 *Let u be a bounded viscosity solution to (1.1) in Ω . Under the assumptions that $0 \leq a(x) \in C^1(\Omega)$, $p \leq q \leq p + \frac{1}{2}$ and (1.3), for any $\Omega' \subset\subset \Omega$, there is a constant C that depends on $n, p, q, \gamma_\infty, \Omega', \Omega, \|a\|_{C^1(\Omega)}, \|u\|_{L^\infty(\Omega)}$ and $\|\Phi\|_{L^\infty(\Omega)}$, such that*

$$|u(x) - u(y)| \leq C|x - y|$$

for all $x, y \in \Omega'$. Here, $\gamma_\infty := \max_{t \in [0, \|u\|_{L^\infty(\Omega)}]} \gamma(t)$.

In order to verify Theorem 1.2, we need to utilize twice the Ishii-Lions methods [31], and to adjust carefully the distance $q - p$ in order to get a contradiction. Combining the above two theorems yields that the bounded viscosity solutions are weak solutions with some explicit conditions.

To show that weak solutions are viscosity solutions, we consider a class of functions satisfying the following comparison principle.

Definition 1.3 Suppose that u is a weak supersolution to (1.1) in $\Omega' \subset \Omega$. If for any weak subsolution v of (1.1) such that $v \leq u$ a.e. in $\partial\Omega'$ there holds that $v \leq u$ a.e. in Ω' , then we say that (u, f) fulfills the comparison principle property (CCP) in Ω' .

Finally, the fact that weak solutions are viscosity solutions can be obtained under the (CCP) condition by a contradiction argument.

Theorem 1.4 *Let u be a lower semicontinuous weak supersolution to (1.1) in Ω . Assume that $f(x, \tau, \xi)$ is uniformly continuous in $\Omega \times \mathbb{R} \times \mathbb{R}^n$. If (CCP) holds true, then u is also a viscosity supersolution to problem (1.1).*

Remark 1.5 When the nonlinear term f only depends on x , and not on u and ∇u , we know from [19] that the weak solutions are locally Lipschitz continuous under the minimal hypotheses that $0 \leq a \in W^{1,d}(\Omega)$, $d > n$, and f belongs to a proper Lorentz space along with $q/p \leq 1 + 1/n - 1/d$, if $n > 2$ and $q/p < p$, if $n = 2$.

This paper is organized as follows. In Sect. 2 we introduce some basic properties of function spaces and concepts of solutions as well as some necessary known results. Section 3 is devoted to proving that viscosity solutions are weak solutions to (1.1), and the reverse implication is showed in Sect. 4, where we also establish the comparison principle for two equations with different nonlinearities. In Sect. 5, we verify that the bounded viscosity solutions of (1.1) are locally Lipschitz continuous, which is the indispensable element of equivalence.

2 Preliminaries

In this section, we summarize some basic properties of the Musielak-Orlicz-Sobolev space $W^{1,H(\cdot)}(\Omega)$. These properties can be found in [11, 28, 45]. In addition, we give different notions of solutions to Eq. (1.1) together with some auxiliary results.

2.1 Function spaces

In the rest of this paper, unless otherwise stated, we always assume that hypothesis (1.2) holds. For all $x \in \Omega$ and $\xi \in \mathbb{R}^n$, we shall use the notation

$$H(x, \xi) := |\xi|^p + a(x)|\xi|^q. \quad (2.1)$$

With abuse of notation, we shall also denote $H(x, \xi)$ when $\xi \in \mathbb{R}$. Observe that the generalized Young function H is a Musielak-Orlicz function fulfilling (Δ_2) and (∇_2) conditions. Here, a Young function H is said to satisfy the Δ_2 -condition provided that there is a constant $c > 0$ such that $H(x, 2t) \leq cH(x, t)$ for every $t \geq 0$. We say that H satisfies the ∇_2 -condition if the Fenchel-Young conjugate H^* of H (see (2.2)), satisfies the Δ_2 -condition.

Let us introduce some important properties, to be used later, of the energy density H given by (2.1). We will keep on denoting $H(x, t) = t^p + a(x)t^q$ for $t \geq 0$, that is, ξ is a non-negative number in (2.1). By the Fenchel-Young conjugate of H , we mean the function

$$H^*(x, t) := \sup_{s \geq 0} \{st - H(x, s)\}. \quad (2.2)$$

It is well known that the equivalence

$$H^*(x, H(x, t)/t) \sim H(x, t) \quad (2.3)$$

holds up to some constants depending on p, q , and moreover the Young inequality

$$st \leq H^*(x, t) + H(x, s) \quad (2.4)$$

holds for all $x \in \Omega, s, t \in [0, +\infty)$.

The Musielak-Orlicz space $L^{H(\cdot)}(\Omega)$ is defined as

$$L^{H(\cdot)}(\Omega) := \{u : \Omega \rightarrow \mathbb{R} \text{ measurable} : \varrho_H(u) < \infty\},$$

and it is endowed with the norm

$$\|u\|_{L^{H(\cdot)}(\Omega)} := \inf \left\{ \lambda > 0 : \varrho_H \left(\frac{u}{\lambda} \right) \leq 1 \right\},$$

where

$$\varrho_H(u) := \int_{\Omega} H(x, u) dx = \int_{\Omega} |u|^p + a(x)|u|^q dx$$

is called the ϱ_H -modular function.

The space $L^{H(\cdot)}(\Omega)$ is a separable, uniformly convex Banach space. From the definitions of the ϱ_H -modular and the norm, we deduce that

$$\min \left\{ \|u\|_{L^{H(\cdot)}(\Omega)}^p, \|u\|_{L^{H(\cdot)}(\Omega)}^q \right\} \leq \varrho_H(u) \leq \max \left\{ \|u\|_{L^{H(\cdot)}(\Omega)}^p, \|u\|_{L^{H(\cdot)}(\Omega)}^q \right\}. \tag{2.5}$$

It follows from (2.5) that

$$\|u_n - u\|_{L^{H(\cdot)}(\Omega)} \rightarrow 0 \iff \varrho_H(u_n - u) \rightarrow 0,$$

which indicates the equivalence of convergence in ϱ_H -modular and in norm. For the space $L^{H^*(\cdot)}(\Omega)$ we know that

$$\begin{aligned} \min \left\{ (\varrho_{H^*}(u))^{\frac{p}{p+1}}, (\varrho_{H^*}(u))^{\frac{q}{q+1}} \right\} &\leq \|u\|_{L^{H^*(\cdot)}(\Omega)} \\ &\leq \max \left\{ (\varrho_{H^*}(u))^{\frac{p}{p+1}}, (\varrho_{H^*}(u))^{\frac{q}{q+1}} \right\}. \end{aligned} \tag{2.6}$$

If $u \in L^{H(\cdot)}(\Omega)$ and $v \in L^{H^*(\cdot)}(\Omega)$, the following Hölder inequality

$$\left| \int_{\Omega} uv \, dx \right| \leq 2 \|u\|_{L^{H(\cdot)}(\Omega)} \|v\|_{L^{H^*(\cdot)}(\Omega)} \tag{2.7}$$

holds.

The Musielak-Orlicz-Sobolev space $W^{1,H(\cdot)}(\Omega)$ is the set of those functions $u \in L^{H(\cdot)}(\Omega)$ satisfying $Du \in L^{H(\cdot)}(\Omega)$. We equip the space $W^{1,H(\cdot)}(\Omega)$ with the norm

$$\|u\|_{W^{1,H(\cdot)}(\Omega)} := \|u\|_{L^{H(\cdot)}(\Omega)} + \|Du\|_{L^{H(\cdot)}(\Omega)}.$$

The space $W^{1,H(\cdot)}(\Omega)$ is a separable and reflexive Banach space. The local space $W_{loc}^{1,H(\cdot)}(\Omega)$ is composed of those functions belonging to $W^{1,H(\cdot)}(\Omega')$ for any sub-domain Ω' compactly involved in Ω . Finally, we denote by $W_0^{1,H(\cdot)}(\Omega)$ the closure of $C_0^\infty(\Omega)$ in $W^{1,H(\cdot)}(\Omega)$. Indeed, the condition (1.2) ensures that the set $C_0^\infty(\Omega)$ is dense in $W_0^{1,H(\cdot)}(\Omega)$ (see [1, 21]).

2.2 Notions of solutions

Set

$$A(x, \xi) := |\xi|^{p-2}\xi + a(x)|\xi|^{q-2}\xi$$

for all $x \in \Omega$ and $\xi \in \mathbb{R}^n$. We now give the definition of diverse type of solutions to problem (1.1). When considering the weak solutions to (1.1), we demand that the source term f fulfills the growth condition (1.3) unless otherwise stated.

Definition 2.1 (weak solution) We say that $u \in W_{loc}^{1,H(\cdot)}(\Omega)$ is a weak supersolution to problem (1.1) if

$$\int_{\Omega} \langle A(x, Du), D\phi \rangle dx \geq \int_{\Omega} f(x, u, Du)\phi dx$$

for each nonnegative function $\phi \in C_0^\infty(\Omega)$. The inequality is reverse for weak subsolution. When $u \in W_{loc}^{1,H(\cdot)}(\Omega)$ is both weak super- and subsolution, we call u a weak solution to (1.1), that is

$$\int_{\Omega} \langle A(x, Du), D\phi \rangle dx = \int_{\Omega} f(x, u, Du)\phi dx$$

for any $\phi \in C_0^\infty(\Omega)$.

Remark 2.2 In the definition above, from (1.3), (2.3) and (2.4) we can see that the integral $\int_{\Omega} f(x, u, Du)\phi dx$ is finite, for any $\phi \in C_0^\infty(\Omega)$.

Let now $\xi, \eta \in \mathbb{R}^n, X \in S^n$ with S^n being the set of symmetric $n \times n$ matrices. We introduce some notations:

$$\begin{aligned} M(x, \xi) &= a(x)|\xi|^{q-2} \left(I + (q-2) \frac{\xi}{|\xi|} \otimes \frac{\xi}{|\xi|} \right), \\ F_1(\xi, X) &= -|\xi|^{p-2} \left(\text{tr}X + (p-2) \left\langle X \frac{\xi}{|\xi|}, \frac{\xi}{|\xi|} \right\rangle \right), \\ F_2(x, \xi, X) &= -a(x)|\xi|^{q-2} \left(\text{tr}X + (q-2) \left\langle X \frac{\xi}{|\xi|}, \frac{\xi}{|\xi|} \right\rangle \right) = -\text{tr}(M(x, \xi)X) \end{aligned}$$

and

$$F_3(x, \xi) = -|\xi|^{q-2} \xi \cdot Da(x),$$

where $\xi \otimes \eta$ denotes an $n \times n$ matrix whose (i, j) entry is $\xi_i \eta_j$, and $\langle \xi, \eta \rangle$ or $\xi \cdot \eta$ stands for the inner product of ξ, η . For a matrix X , we set the matrix norm $\|X\| := \sup_{|\xi| \leq 1} \{ |X\xi| \}$. In order to define the viscosity solutions of problem (1.1), we let $a \in C^1(\Omega)$ and easily check that

$$\begin{aligned} & -\text{div}(|Du|^{p-2} Du + a(x)|Du|^{q-2} Du) \\ &= F_1(Du, D^2u) + F_2(x, Du, D^2u) + F_3(x, Du) \\ &=: F(x, Du, D^2u). \end{aligned} \tag{2.8}$$

We now recall the notion of semi-jets. The subset of $u : \Omega \rightarrow \mathbb{R}$ at x is given by letting $(\eta, X) \in J^{2,-}u(x)$ if

$$u(y) \geq u(x) + \eta \cdot (y - x) + \frac{1}{2} \langle X(y - x), (y - x) \rangle + o(|y - x|^2)$$

as $y \rightarrow x$. The closure of a subset is defined by $(\eta, X) \in \overline{J^{2,-}}u(x)$ if there exists a sequence $(\eta_j, X_j) \in J^{2,-}u(x_j)$ such that $(x_j, u(x_j), \eta_j, X_j) \rightarrow (x, r, \eta, X)$ with some $r \in \mathbb{R}$ as $j \rightarrow \infty$. Obviously, $r = u(x)$ if u is continuous. The superjet $J^{2,+}$ and its closure $\overline{J^{2,+}}$ are defined by a similar way but the inequality above needs to be converse.

Definition 2.3 (viscosity solution) A lower semicontinuous function $u : \Omega \rightarrow (-\infty, \infty)$ is a viscosity supersolution to problem (1.1) in Ω , if $(\eta, X) \in J^{2,-}u(x)$ with $x \in \Omega$ and $\eta \neq 0$ implies that

$$F(x, \eta, X) \geq f(x, u(x), \eta).$$

An upper semicontinuous function $u : \Omega \rightarrow (-\infty, \infty)$ is a viscosity subsolution to problem (1.1) in Ω , if for each $(\eta, X) \in J^{2,+}u(x)$ with $x \in \Omega$ and $\eta \neq 0$ there holds that

$$F(x, \eta, X) \leq f(x, u(x), \eta).$$

A function u is called viscosity solution to (1.1) if and only if it is viscosity super- and subsolution.

Remark 2.4 The preceding concept of viscosity solutions is equivalently given by the jet-closures or test functions. For instance, the following conditions are equivalent:

- (1) A function u is a viscosity supersolution to (1.1) in Ω ;
- (2) If $(\eta, X) \in \overline{J^{2,-}}u(x)$ with $x \in \Omega$ and $\eta \neq 0$, then $F(x, \eta, X) \geq f(x, u(x), \eta)$;
- (3) If $\varphi \in C^2(\Omega)$ touches u from below at x , that is, $\varphi(x) = u(x)$, $\varphi(y) \leq u(y)$ and moreover $D\varphi(x) \neq 0$, then we have $F(x, D\varphi(x), D^2\varphi(x)) \geq f(x, u(x), D\varphi(x))$.

In the case $2 \leq p \leq q$, we can remove the requirement that $\eta \neq 0$ or $D\varphi(x) \neq 0$.

2.3 Inf-convolution

We now give the definition of infimal convolution together with some properties. Define the inf-convolution as

$$u_\varepsilon(x) = \inf_{y \in \Omega} \left\{ u(y) + \frac{|x - y|^s}{s\varepsilon^{s-1}} \right\},$$

where $\varepsilon > 0$ and $s \geq \max \left\{ 2, \frac{p}{p-1} \right\}$ is a constant to be fixed by the growth powers in Eq. (1.1). Indeed, when $2 \leq p \leq q$, $s = 2$; when $1 < p \leq q < 2$, $s > \max \left\{ \frac{p}{p-1}, \frac{q}{q-1} \right\} = \frac{p}{p-1}$; when $1 < p < 2 \leq q$, $s > \max \left\{ \frac{p}{p-1}, 2 \right\} = \frac{p}{p-1}$.

The following well-known properties of the inf-convolution u_ε can be found in several references, such as [32, 49].

Proposition 2.5 Suppose that $u : \Omega \rightarrow \mathbb{R}$ is a bounded and lower semicontinuous function. Then the inf-convolution u_ε satisfies the following properties:

- (1) $u_\varepsilon \leq u$ in Ω and $u_\varepsilon \rightarrow u$ locally uniformly as $\varepsilon \rightarrow 0$.
- (2) There is $r(\varepsilon) > 0$ such that

$$u_\varepsilon(x) = \inf_{y \in B_{r(\varepsilon)}(x) \cap \Omega} \left\{ u(y) + \frac{|x - y|^s}{s\varepsilon^{s-1}} \right\}$$

with $r(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$. In particular, if $x \in \Omega_{r(\varepsilon)} := \{x \in \Omega : \text{dist}(x, \partial\Omega) > r(\varepsilon)\}$, then there exists a point $x_\varepsilon \in B_{r(\varepsilon)}(x)$ fulfilling

$$u_\varepsilon(x) = u(x_\varepsilon) + \frac{|x - x_\varepsilon|^s}{s\varepsilon^{s-1}}.$$

- (3) The function u_ε is semi-concave in $\Omega_{r(\varepsilon)}$, that is, we can find a constant C , depending only on u, s and ε , such that the function $x \mapsto u_\varepsilon(x) - C|x|^2$ is concave.
- (4) If $(\eta, X) \in J^{2,-}u_\varepsilon(x)$ with $x \in \Omega_{r(\varepsilon)}$, then we have

$$\eta = \frac{|x - x_\varepsilon|^{s-2}(x - x_\varepsilon)}{\varepsilon^{s-1}} \quad \text{and} \quad X \leq \frac{s-1}{\varepsilon} |\eta|^{\frac{s-2}{s-1}} I.$$

3 Viscosity solutions are weak solutions

In this part, we are going to make use of the inf-convolution approximation technique to show that the locally Lipschitz continuous viscosity solutions are weak solutions to (1.1). Furthermore, the Lipschitz continuity of viscosity solutions can be established precisely; this proof is postponed to Sect. 5. We therefore draw a conclusion that viscosity solutions are weak solutions under some suitable conditions. We will only discuss viscosity and weak supersolutions below. The case of subsolutions is similar.

We begin by stating that if u is a (locally Lipschitz) viscosity supersolution to (1.1) in Ω , then its inf-convolution u_ε is also a viscosity supersolution of such equation (whose form may be slightly changed) in a shrinking domain. From the following lemma, we can find that owing to the presence of the modulating coefficient $a(x)$, there will exist an error term $E(\varepsilon)$ on the right-hand side of the equation. In what follows, $X \leq Y$ ($X, Y \in \mathcal{S}^n$) means that $\langle (X - Y)\xi, \xi \rangle \leq 0$ for any $\xi \in \mathbb{R}^n$. We denote by C a generic constant, which may vary from line to line. If necessary, relevant dependencies on parameters will be emphasised using parentheses.

Lemma 3.1 *Assume that $0 < a(x) \in C^1(\Omega)$ and $f(x, \tau, \xi)$ is continuous in $\Omega \times \mathbb{R} \times \mathbb{R}^n$ and decreasing with respect to τ . Let u be a viscosity supersolution with local Lipschitz to (1.1) in Ω . Then if $(\eta, X) \in J^{2,-}u_\varepsilon(x)$ with $\eta \neq 0$ and $x \in \Omega_{r(\varepsilon)}$, there holds that*

$$F(x, \eta, X) \geq f_\varepsilon(x, u_\varepsilon(x), \eta) + E(\varepsilon), \tag{3.1}$$

where

$$f_\varepsilon(x, \tau, \xi) := \inf_{y \in B_{r(\varepsilon)}(x)} f(y, \tau, \xi)$$

and $E(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$.

Remark 3.2 It is worth mentioning that the requirement $a(x) > 0$ is just needed in this lemma for the technical reason.

Proof Let $x^o \in \Omega_{r(\varepsilon)}$ and $(\eta, X) \in J^{2,-}u_\varepsilon(x^o)$ with $\eta \neq 0$. Then via the properties of the inf-convolution u_ε in Proposition 2.5, we have

$$u_\varepsilon(x^o) = u(x_\varepsilon^o) + \frac{|x^o - x_\varepsilon^o|^s}{s\varepsilon^{s-1}}$$

with $x_\varepsilon^o \in B_{r(\varepsilon)}(x^o)$, and moreover $\eta = \frac{|x^o - x_\varepsilon^o|^{s-2}}{\varepsilon^{s-1}}(x^o - x_\varepsilon^o)$. There is a function $\varphi \in C^2(\Omega)$ such that it touches u_ε from below at x^o and $D\varphi(x^o) = \eta, D^2\varphi(x^o) = X$. Hence by the definition of inf-convolution u_ε we can see that

$$0 \leq u_\varepsilon(x) - \varphi(x) \leq u(y) + \frac{|x - y|^s}{s\varepsilon^{s-1}} - \varphi(x)$$

for any $x, y \in \Omega_{r(\varepsilon)}$. Notice also that

$$u(x_\varepsilon^o) + \frac{|x^o - x_\varepsilon^o|^s}{s\varepsilon^{s-1}} - \varphi(x^o) = u_\varepsilon(x^o) - \varphi(x^o) = 0.$$

Then the function $-u(y) + \varphi(x) - \frac{|y-x|^s}{s\varepsilon^{s-1}}$ attains the maximum at (x_ε^o, x^o) . Therefore, by applying the maximum principle for semicontinuous functions (also known as the theorem of sums) in [16], we can find $Y, Z \in \mathcal{S}^n$ satisfying

$$(-D_y\psi(x_\varepsilon^o, x^o), -Y) \in \bar{J}^{2,-}u(x_\varepsilon^o), \quad (D_x\psi(x_\varepsilon^o, x^o), -Z) \in \bar{J}^{2,+}\varphi(x^o)$$

and

$$\begin{pmatrix} Y \\ -Z \end{pmatrix} \leq D^2\psi(x_\varepsilon^o, x^o) + \varepsilon^{1-s}(D^2\psi(x_\varepsilon^o, x^o))^2,$$

where $\psi(y, x) := \frac{|y-x|^s}{s\varepsilon^{s-1}}$ and

$$D^2\psi(x_\varepsilon^o, x^o) = \begin{pmatrix} D_{yy}\psi(x_\varepsilon^o, x^o) & D_{yx}\psi(x_\varepsilon^o, x^o) \\ D_{xy}\psi(x_\varepsilon^o, x^o) & D_{xx}\psi(x_\varepsilon^o, x^o) \end{pmatrix}.$$

Via direct computation, we obtain

$$-D_y\psi(x_\varepsilon^o, x^o) = \eta = D_x\psi(x_\varepsilon^o, x^o)$$

and

$$B := D_{yy}\psi(x_\varepsilon^o, x^o) = \varepsilon^{1-s}|x_\varepsilon^o - x^o|^{s-4}[|x_\varepsilon^o - x^o|^2I + (s-2)(x_\varepsilon^o - x^o) \otimes (x_\varepsilon^o - x^o)].$$

Furthermore,

$$\begin{pmatrix} Y & \\ & -Z \end{pmatrix} \leq \begin{pmatrix} B & -B \\ -B & B \end{pmatrix} + 2\varepsilon^{1-s} \begin{pmatrix} B^2 & -B^2 \\ -B^2 & B^2 \end{pmatrix}. \tag{3.2}$$

It is straightforward to derive that $Y \leq Z$ and moreover $Z \leq -X$. Since u is a viscosity supersolution to (1.1) and f is continuous in all variables, then

$$\begin{aligned} f(x_\varepsilon^o, u(x_\varepsilon^o), \eta) &\leq F(x_\varepsilon^o, \eta, -Y) \\ &= F_1(\eta, Z) - F_1(\eta, Y) - F_1(\eta, Z) + F_2(x^o, \eta, Z) - F_2(x_\varepsilon^o, \eta, Y) \\ &\quad - F_2(x^o, \eta, Z) + F_3(x^o, \eta) + F_3(x_\varepsilon^o, \eta) - F_3(x^o, \eta) \\ &\leq F(x^o, \eta, -Z) + F_2(x^o, \eta, Z) - F(x_\varepsilon^o, \eta, Y) + F_3(x_\varepsilon^o, \eta) - F_3(x^o, \eta) \end{aligned}$$

with the notations given in (2.8), where in the last line we have used the decreasing property of $F_1(\xi, X)$ in the X -variable, that is, $Y \leq Z$ implies $F_1(\eta, Z) \leq F_1(\eta, Y)$. Next, in view of $a \in C^1(\Omega)$ we estimate

$$\begin{aligned} F_3(x_\varepsilon^o, \eta) - F_3(x^o, \eta) &= |\eta|^{q-2} \eta \cdot Da(x^o) - |\eta|^{q-2} \eta \cdot Da(x_\varepsilon^o) \\ &\leq |\eta|^{q-1} |Da(x^o) - Da(x_\varepsilon^o)| \\ &\leq |\eta|^{q-1} \omega(r(\varepsilon)), \end{aligned} \tag{3.3}$$

where $\omega(\cdot)$ represents the modulus of continuity of Da . We finally evaluate the term $F_2(x^o, \eta, Z) - F_2(x_\varepsilon^o, \eta, Y)$ in a similar way to address $F_2(y_j, \eta_j, Y_j) - F_2(x_j, \eta_j, X_j)$ in [23, Proposition 5.1]. Observe that it follows from (3.2) that

$$\langle \xi, Y\xi \rangle - \langle \zeta, Z\xi \rangle \leq \varepsilon^{1-s} \left[(s-1)|x_\varepsilon^o - x^o|^{s-2} + 2(s-1)^2|x_\varepsilon^o - x^o|^{2(s-2)} \right] |\xi - \zeta|^2 \tag{3.4}$$

with $\xi, \zeta \in \mathbb{R}^n$. We can easily verify the matrix $M(x, \xi) \geq 0$ (positive semi-definite) as $a(x) \geq 0$ so that it has square root denoted by $M^{\frac{1}{2}}(x, \xi)$. Additionally, by $M_l^{\frac{1}{2}}(x, \xi)$ we mean the l -th column of $M^{\frac{1}{2}}(x, \xi)$. Then employing (3.4) and decomposition of matrix yields that

$$\begin{aligned} &F_2(x^o, \eta, Z) - F_2(x_\varepsilon^o, \eta, Y) \\ &= \text{tr} \left(M^{\frac{1}{2}}(x^o, \eta) M^{\frac{1}{2}}(x^o, \eta) Z \right) - \text{tr} \left(M^{\frac{1}{2}}(x_\varepsilon^o, \eta) M^{\frac{1}{2}}(x_\varepsilon^o, \eta) Y \right) \\ &= \sum_{l=1}^n \left\langle M_l^{\frac{1}{2}}(x^o, \eta), Z M_l^{\frac{1}{2}}(x^o, \eta) \right\rangle - \sum_{l=1}^n \left\langle M_l^{\frac{1}{2}}(x_\varepsilon^o, \eta), Y M_l^{\frac{1}{2}}(x_\varepsilon^o, \eta) \right\rangle \\ &\leq C\varepsilon^{1-s} |x_\varepsilon^o - x^o|^{s-2} \left\| M^{\frac{1}{2}}(x^o, \eta) - M^{\frac{1}{2}}(x_\varepsilon^o, \eta) \right\|_2^2 \\ &\leq \frac{C\varepsilon^{1-s} |x_\varepsilon^o - x^o|^{s-2}}{\left(\lambda_{\min} \left(M^{\frac{1}{2}}(x^o, \eta) \right) + \lambda_{\min} \left(M^{\frac{1}{2}}(x_\varepsilon^o, \eta) \right) \right)^2} \|M(x^o, \eta) - M(x_\varepsilon^o, \eta)\|_2^2 \end{aligned}$$

$$\leq \frac{C\varepsilon^{1-s}|x_\varepsilon^o - x^o|^{s-2}|\eta|^{2(q-2)}|a(x^o) - a(x_\varepsilon^o)|^2}{|\eta|^{q-2}(\sqrt{a(x^o)} + \sqrt{a(x_\varepsilon^o)})^2},$$

where $\lambda_{\min}(M)$ stands for the smallest eigenvalue of the matrix M . More details of this display can be found in [23, Proposition 5.1].

On the other hand, from the local Lipschitz continuity of u ,

$$\frac{|x^o - x_\varepsilon^o|^s}{s\varepsilon^{s-1}} = u_\varepsilon(x^o) - u(x_\varepsilon^o) < u(x^o) - u(x_\varepsilon^o) \leq C|x_\varepsilon^o - x^o|,$$

we have

$$|\eta| = \frac{|x^o - x_\varepsilon^o|^{s-1}}{\varepsilon^{s-1}} \leq C.$$

By means of $a(x) \in C^1(\Omega)$, we proceed to treat

$$\begin{aligned} F_2(x^o, \eta, Z) - F_2(x_\varepsilon^o, \eta, Y) &\leq \frac{C\varepsilon^{1-s}|x_\varepsilon^o - x^o|^{s-2}|\eta|^{q-2}|x^o - x_\varepsilon^o|^2}{(\sqrt{a(x^o)} + \sqrt{a(x_\varepsilon^o)})^2} \\ &= \frac{C|\eta|^{q-1}|x^o - x_\varepsilon^o|}{(\sqrt{a(x^o)} + \sqrt{a(x_\varepsilon^o)})^2} \\ &\leq \frac{Cr(\varepsilon)}{(\sqrt{a(x^o)} + \sqrt{a(x_\varepsilon^o)})^2}. \end{aligned}$$

Here we remark that if $a(x^o) = 0$, then the quantity $\frac{|\eta|^{q-1}|x^o - x_\varepsilon^o|}{a(x_\varepsilon^o)}$ does not necessarily go to 0 as $\varepsilon \rightarrow 0$, so the condition $a(x) > 0$ is required. Besides, by the boundedness of η the inequality (3.3) becomes

$$F_3(x_\varepsilon^o, \eta) - F_3(x^o, \eta) \leq C\omega(r(\varepsilon)).$$

Since $f(x, \tau, \xi)$ is decreasing in τ and $u(x_\varepsilon^o) \leq u_\varepsilon(x^o)$,

$$f(x_\varepsilon^o, u(x_\varepsilon^o), \eta) \geq f(x_\varepsilon^o, u_\varepsilon(x^o), \eta) \geq \inf_{y \in B_{r(\varepsilon)}(x^o)} f(y, u_\varepsilon(x^o), \eta).$$

Now define

$$E(\varepsilon) := -C\omega(r(\varepsilon)) - \frac{Cr(\varepsilon)}{(\sqrt{a(x^o)} + \sqrt{a(x_\varepsilon^o)})^2}.$$

Consequently, we get

$$f_\varepsilon(x^o, u_\varepsilon(x^o), \eta) + E(\varepsilon) \leq F(x^o, \eta, X)$$

with $E(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$. Here we have employed $F(x^o, \eta, X) \geq F(x^o, \eta, -Z)$ by $Z \leq -X$ and the error term $E(\varepsilon)$ depends on ε, n, p, q, a and the Lipschitz constant of u . The proof is now completed. \square

Based on the preceding lemma, we further demonstrate that when u is a viscosity supersolution to (1.1), then the inf-convolution u_ε is a weak supersolution of this equation up to a certain error term. Let $\Delta_\infty u = \langle Du, D^2u Du \rangle$ below.

Lemma 3.3 *Suppose that the assumptions on a, f in Lemma 3.1 are in force, and that $f(x, \tau, 0) \leq 0$ for all $(x, \tau) \in \Omega \times \mathbb{R}$. Let u be a locally Lipschitz continuous viscosity supersolution to (1.1). Then, for each nonnegative function $\varphi \in W_0^{1,H(\cdot)}(\Omega_{r(\varepsilon)})$, it holds that*

$$\int_{\Omega_{r(\varepsilon)}} \varphi f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) dx + E(\varepsilon) \int_{\{Du_\varepsilon \neq 0\}} \varphi dx \leq \int_{\Omega_{r(\varepsilon)}} \langle A(x, Du_\varepsilon), D\varphi \rangle dx,$$

where $E(\varepsilon) \rightarrow 0$, which is from Lemma 3.1, as $\varepsilon \rightarrow 0$.

Proof It suffices to consider the nonnegative function $\varphi \in C_0^\infty(\Omega_{r(\varepsilon)})$, because the function space $C_0^\infty(\Omega_{r(\varepsilon)})$ is dense in $W_0^{1,H(\cdot)}(\Omega_{r(\varepsilon)})$.

Owing to u_ε being semi-concave, we can see by Proposition 2.5 that

$$h(x) := u_\varepsilon(x) - C(s, \varepsilon, u)|x|^2$$

is concave in $\Omega_{r(\varepsilon)}$. Let $\{h_j\}_j$ be a sequence of smooth concave functions, obtained from standard mollification, such that

$$(h_j, Dh_j, D^2h_j) \rightarrow (h, Dh, D^2h) \text{ a.e. in } \Omega_{r(\varepsilon)}.$$

Additionally, we define

$$u_{\varepsilon,j} = h_j + C(s, \varepsilon, u)|x|^2$$

and let $\delta \in (0, 1)$. In this proof, the condition $0 \leq a(x) \in C^{0,1}(\Omega)$ is sufficient except applying Lemma 3.1. Now denote the standard mollification of a as a_j .

Case 1. $1 < p \leq q < 2$. In this singular case, we first regularize the equation by adding a small $\delta > 0$ as follows, and eventually pass to the limit as $\delta \rightarrow 0$. Recalling that $u_{\varepsilon,j}$ and a_j are smooth, we can calculate by integration by parts

$$\begin{aligned} & \int_{\Omega_{r(\varepsilon)}} -\varphi \operatorname{div} \left[(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} Du_{\varepsilon,j} + a_j(x)(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} Du_{\varepsilon,j} \right] dx \\ &= \int_{\Omega_{r(\varepsilon)}} \left\langle (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} Du_{\varepsilon,j} + a_j(x)(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} Du_{\varepsilon,j}, D\varphi \right\rangle dx. \end{aligned} \tag{3.5}$$

We are ready to consider the limit as $j \rightarrow \infty$, and claim that

$$\begin{aligned}
 & - \int_{\Omega_{r(\varepsilon)}} \varphi \operatorname{div} \left[(|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}} Du_\varepsilon + a(x)(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon \right] dx \\
 & \leq \int_{\Omega_{r(\varepsilon)}} \left\langle (|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}} Du_\varepsilon + a(x)(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon, D\varphi \right\rangle dx. \tag{3.6}
 \end{aligned}$$

It follows, from the Lipschitz continuity of u_ε and a , that for some constant $M > 0$

$$\|Du_{\varepsilon,j}\|_{L^\infty(\operatorname{supp}'\varphi)}, \|a_j\|_{L^\infty(\operatorname{supp}'\varphi)}, \|Da_j\|_{L^\infty(\operatorname{supp}'\varphi)} \leq M, \quad j = 1, 2, 3, \dots$$

Thus we could apply the Lebesgue dominated convergence theorem to the integral at the right-hand side of (3.5). On the other hand, via direct computation,

$$\begin{aligned}
 & - \int_{\Omega_{r(\varepsilon)}} \varphi \operatorname{div} \left[(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} Du_{\varepsilon,j} + a_j(x)(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} Du_{\varepsilon,j} \right] dx \\
 & = - \int_{\Omega_{r(\varepsilon)}} \varphi (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} \left(\Delta u_{\varepsilon,j} + \frac{p-2}{|Du_{\varepsilon,j}|^2 + \delta} \Delta_\infty u_{\varepsilon,j} \right) dx \\
 & \quad - \int_{\Omega_{r(\varepsilon)}} \varphi a_j (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} \left(\Delta u_{\varepsilon,j} + \frac{q-2}{|Du_{\varepsilon,j}|^2 + \delta} \Delta_\infty u_{\varepsilon,j} \right) dx \\
 & \quad - \int_{\Omega_{r(\varepsilon)}} \varphi (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} Du_{\varepsilon,j} \cdot Da_j dx. \tag{3.7}
 \end{aligned}$$

Obviously, by the dominated convergence theorem, when $j \rightarrow \infty$,

$$\int_{\Omega_{r(\varepsilon)}} \varphi (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} Du_{\varepsilon,j} \cdot Da_j dx \rightarrow \int_{\Omega_{r(\varepsilon)}} \varphi (|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon \cdot Da dx.$$

Note that h_j is concave. Then we have $D^2u_{\varepsilon,j} \leq C(s, \varepsilon, u)I$. For $Du_{\varepsilon,j} \neq 0$, we arrive at

$$(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} \left(\Delta u_{\varepsilon,j} + \frac{p-2}{|Du_{\varepsilon,j}|^2 + \delta} \Delta_\infty u_{\varepsilon,j} \right) \leq C(s, \varepsilon, u) \delta^{\frac{p-2}{2}} (2n + p - 2)$$

and

$$\begin{aligned}
 & a_j(x)(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{q-2}{2}} \left(\Delta u_{\varepsilon,j} + \frac{q-2}{|Du_{\varepsilon,j}|^2 + \delta} \Delta_\infty u_{\varepsilon,j} \right) \\
 & \leq MC(s, \varepsilon, u) \delta^{\frac{q-2}{2}} (2n + q - 2).
 \end{aligned}$$

For $Du_{\varepsilon,j} = 0$, the case becomes easier. Therefore, we can apply Fatou’s lemma to the display (3.7), and further justify (3.6).

Next, we shall let $\delta \rightarrow 0$ in the inequality (3.6). It follows from the dominated convergence theorem that, when δ goes to 0,

$$\begin{aligned} & \int_{\Omega_{r(\varepsilon)}} \left\langle (|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}} Du_\varepsilon + a(x)(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon, D\varphi \right\rangle dx \\ & \rightarrow \int_{\Omega_{r(\varepsilon)}} \left\langle |Du_\varepsilon|^{p-2} Du_\varepsilon + a(x)|Du_\varepsilon|^{q-2} Du_\varepsilon, D\varphi \right\rangle dx. \end{aligned}$$

Besides, we show that the integrand at the left-hand side of (3.6) is bounded from below, which can justify the use of Fatou’s lemma. If $Du_\varepsilon = 0$, this follows immediately from the inequality (see Proposition 2.5)

$$D^2u_\varepsilon \leq \frac{s-1}{\varepsilon} |Du_\varepsilon|^{\frac{s-2}{s-1}} I.$$

In other words, since $s > 2$, then D^2u_ε is negative semi-definite when $Du_\varepsilon = 0$. If $Du_\varepsilon \neq 0$, we will find that

$$\begin{aligned} & -(|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}} \left(\Delta u_\varepsilon + \frac{p-2}{|Du_\varepsilon|^2 + \delta} \Delta_\infty u_\varepsilon \right) \\ & \geq -\frac{(|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}}}{|Du_\varepsilon|^2 + \delta} \frac{s-1}{\varepsilon} \left(|Du_\varepsilon|^{\frac{s-2}{s-1}+2} (n+p-2) + \delta n |Du_\varepsilon|^{\frac{s-2}{s-1}} \right) \\ & \geq -|Du_\varepsilon|^{\frac{s-2}{s-1}+p-2} \frac{s-1}{\varepsilon} (2n+p-2) \\ & \geq -\|Du_\varepsilon\|_{L^\infty(\Omega_{r(\varepsilon)})}^{\frac{s-2}{s-1}+p-2} \frac{s-1}{\varepsilon} (2n+p-2), \end{aligned}$$

where in the last inequality we need to recall the local Lipschitz continuity of u_ε and $s > \frac{p}{p-1}$. Likewise,

$$\begin{aligned} & -a(x)(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} \left(\Delta u_\varepsilon + \frac{q-2}{|Du_\varepsilon|^2 + \delta} \Delta_\infty u_\varepsilon \right) \\ & \geq -\|a\|_{L^\infty(\Omega_{r(\varepsilon)})} \|Du_\varepsilon\|_{L^\infty(\Omega_{r(\varepsilon)})}^{\frac{s-2}{s-1}+q-2} \frac{s-1}{\varepsilon} (2n+q-2). \end{aligned}$$

Here we note $s > \frac{p}{p-1} \geq \frac{q}{q-1}$. Moreover, we have

$$\left| (|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon \cdot Da \right| \leq \|Da\|_{L^\infty(\Omega_{r(\varepsilon)})} \|Du_\varepsilon\|_{L^\infty(\Omega_{r(\varepsilon)})}^{q-1}.$$

On the other hand, we know that $(Du_\varepsilon(x), D^2u_\varepsilon(x)) \in J^{2,-}u_\varepsilon(x)$ for almost every $x \in \Omega_{r(\varepsilon)}$. Then by means of Lemma 3.1 we deduce

$$F(x, Du_\varepsilon(x), D^2u_\varepsilon(x)) \geq f_\varepsilon(x, u_\varepsilon(x), Du_\varepsilon(x)) + E(\varepsilon) \tag{3.8}$$

in $\{x \in \Omega_{r(\varepsilon)} : Du_\varepsilon \neq 0\}$. As a consequence, exploiting Fatou’s lemma along with observing that D^2u_ε is negative semi-definite when $Du_\varepsilon = 0$, we derive from (3.6) that

$$\begin{aligned}
 & \int_{\Omega_{r(\varepsilon)}} \left\langle |Du_\varepsilon|^{p-2} Du_\varepsilon + a(x)|Du_\varepsilon|^{q-2} Du_\varepsilon, D\varphi \right\rangle dx \\
 & \geq \liminf_{\delta \rightarrow 0} \left[\int_{\{Du_\varepsilon \neq 0\}} + \int_{\{Du_\varepsilon = 0\}} -\varphi(|Du_\varepsilon|^2 + \delta)^{\frac{p-2}{2}} \left(\Delta u_\varepsilon + \frac{p-2}{|Du_\varepsilon|^2 + \delta} \Delta_\infty u_\varepsilon \right) dx \right] \\
 & \quad + \liminf_{\delta \rightarrow 0} \left[\int_{\{Du_\varepsilon \neq 0\}} + \int_{\{Du_\varepsilon = 0\}} -\varphi a(x)(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} \left(\Delta u_\varepsilon + \frac{q-2}{|Du_\varepsilon|^2 + \delta} \Delta_\infty u_\varepsilon \right) dx \right] \\
 & \quad + \liminf_{\delta \rightarrow 0} \left[\int_{\{Du_\varepsilon \neq 0\}} + \int_{\{Du_\varepsilon = 0\}} -\varphi(|Du_\varepsilon|^2 + \delta)^{\frac{q-2}{2}} Du_\varepsilon \cdot Da \, dx \right] \\
 & \geq \int_{\{Du_\varepsilon \neq 0\}} -\varphi |Du_\varepsilon|^{p-2} \left(\Delta u_\varepsilon + \frac{p-2}{|Du_\varepsilon|^2} \Delta_\infty u_\varepsilon \right) dx \\
 & \quad + \int_{\{Du_\varepsilon \neq 0\}} -\varphi a(x) |Du_\varepsilon|^{q-2} \left(\Delta u_\varepsilon + \frac{q-2}{|Du_\varepsilon|^2} \Delta_\infty u_\varepsilon \right) dx \\
 & \quad + \int_{\{Du_\varepsilon \neq 0\}} -\varphi |Du_\varepsilon|^{q-2} Du_\varepsilon \cdot Da \, dx \\
 & = \int_{\{Du_\varepsilon \neq 0\}} \varphi F(x, Du_\varepsilon, D^2u_\varepsilon) \, dx \\
 & \geq \int_{\{Du_\varepsilon \neq 0\}} \varphi f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) \, dx + E(\varepsilon) \int_{\{Du_\varepsilon \neq 0\}} \varphi \, dx \\
 & \geq \int_{\Omega_{r(\varepsilon)}} \varphi f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) \, dx + E(\varepsilon) \int_{\{Du_\varepsilon \neq 0\}} \varphi \, dx, \tag{3.9}
 \end{aligned}$$

where in the penultimate line we employed the inequality (3.8), and in the last line we used the assumption that $f(x, \tau, 0) \leq 0$.

Case 2. $1 < p < 2 \leq q$. We need to substitute the identity (3.5) with

$$\begin{aligned}
 & \int_{\Omega_{r(\varepsilon)}} -\varphi \operatorname{div} \left[(|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} Du_{\varepsilon,j} + a_j(x) |Du_{\varepsilon,j}|^{q-2} Du_{\varepsilon,j} \right] dx \\
 & = \int_{\Omega_{r(\varepsilon)}} \left\langle (|Du_{\varepsilon,j}|^2 + \delta)^{\frac{p-2}{2}} Du_{\varepsilon,j} + a_j(x) |Du_{\varepsilon,j}|^{q-2} Du_{\varepsilon,j}, D\varphi \right\rangle dx,
 \end{aligned}$$

since the q -growth term does not have singularity so that it is not necessarily regularized as the p -growth term. Then the subsequent processes are analogous to Case 1.

Case 3. $2 \leq p \leq q$. In this non-singular scenario, we replace the display (3.5) by

$$\begin{aligned}
 & \int_{\Omega_{r(\varepsilon)}} -\varphi \operatorname{div} \left[|Du_{\varepsilon,j}|^{p-2} Du_{\varepsilon,j} + a_j(x) |Du_{\varepsilon,j}|^{q-2} Du_{\varepsilon,j} \right] dx \\
 & = \int_{\Omega_{r(\varepsilon)}} \left\langle |Du_{\varepsilon,j}|^{p-2} Du_{\varepsilon,j} + a_j(x) |Du_{\varepsilon,j}|^{q-2} Du_{\varepsilon,j}, D\varphi \right\rangle dx.
 \end{aligned}$$

The other procedures are also similar to Case 1, even more straightforward due to the absence of δ . All in all, we finally deduce the desired result. □

In the previous lemma, in order to get the inequality (3.9), $f(x, \tau, 0)$ is *a priori* assumed to be non-positive. The forthcoming lemma states this hypotheses exactly can be realized when $Du_\varepsilon = 0$.

Lemma 3.4 *Let u be a bounded viscosity supersolution to (1.1) in Ω . Let also $0 \leq a(x) \in C^1(\Omega)$ and the function $f(x, \tau, \xi)$ be continuous in all variables. Whenever $Du_\varepsilon(\hat{x}) = 0$ for some $\hat{x} \in \Omega_{r(\varepsilon)}$, we have*

$$f_\varepsilon(\hat{x}, u_\varepsilon(\hat{x}), Du_\varepsilon(\hat{x})) \leq 0.$$

Proof We have known by [32, Lemma 4.3] that if $Du_\varepsilon(\hat{x}) = 0$, then

$$u_\varepsilon(\hat{x}) = u(\hat{x}).$$

From the definition of inf-convolution,

$$u(\hat{x}) \leq u(y) + \frac{|\hat{x} - y|^s}{s\varepsilon^{s-1}} \text{ for all } y \in \Omega.$$

Now introduce an auxiliary function

$$\phi(y) = u(\hat{x}) - \frac{|\hat{x} - y|^s}{s\varepsilon^{s-1}} \text{ with } y \in \Omega,$$

where $s = 2$ if $p \geq 2$, and $s > \max\left\{\frac{p}{p-1}, \frac{q}{q-1}\right\} = \frac{p}{p-1}$ if $1 < p < 2$. We can apparently see that

$$\phi \in C^2(\Omega), D\phi(\hat{x}) = 0, D\phi(y) \neq 0 \text{ for } y \neq \hat{x},$$

and ϕ touches u from below at \hat{x} . For the case $s > 2$ (i.e., $1 < p < 2$), we next evaluate some important quantities,

$$\begin{aligned} D\phi &= \varepsilon^{1-s}|\hat{x} - y|^{s-2}(\hat{x} - y), \\ D^2\phi &= -\varepsilon^{1-s}|\hat{x} - y|^{s-2} \left(I + (s-2) \frac{\hat{x} - y}{|\hat{x} - y|} \otimes \frac{\hat{x} - y}{|\hat{x} - y|} \right) \end{aligned}$$

and

$$\begin{aligned} \text{tr} D^2\phi &= -(n + s - 2)\varepsilon^{1-s}|\hat{x} - y|^{s-2}, \\ \left\langle \frac{D\phi}{|D\phi|}, D^2\phi \frac{D\phi}{|D\phi|} \right\rangle &= -(s - 1)\varepsilon^{1-s}|\hat{x} - y|^{s-2}. \end{aligned}$$

Combining these quantities leads to

$$\begin{aligned} &\text{div}(|D\phi|^{p-2}D\phi + a(y)|D\phi|^{p-2}D\phi) \\ &= \Delta_p\phi + a(y)\Delta_q\phi + |D\phi|^{p-2}D\phi \cdot Da(y) \end{aligned}$$

$$\begin{aligned}
 &= -(n + (p - 1)(s - 1) - 1)\varepsilon^{(1-s)(p-1)}|\hat{x} - y|^{(s-1)(p-1)-1} \\
 &\quad - a(y)(n + (q - 1)(s - 1) - 1)\varepsilon^{(1-s)(q-1)}|\hat{x} - y|^{(s-1)(q-1)-1} \\
 &\quad + \varepsilon^{(1-s)(q-1)}|\hat{x} - y|^{(s-1)(q-1)-1}(\hat{x} - y) \cdot Da(y).
 \end{aligned}$$

Observe that $s > \max \left\{ \frac{p}{p-1}, \frac{q}{q-1} \right\}$. It implies that

$$(s - 1)(p - 1) - 1 > 0 \quad \text{and} \quad (s - 1)(q - 1) - 1 > 0.$$

Thereby,

$$\lim_{r \rightarrow 0} \sup_{y \in B_r(\hat{x}) \setminus \{\hat{x}\}} (-\operatorname{div}(|D\phi|^{p-2}D\phi + a(y)|D\phi|^{p-2}D\phi)) = 0.$$

As for $s = 2$ (i.e., $p \geq 2$), the previous identity is valid obviously. Because u is a viscosity supersolution of (1.1), there holds that

$$\lim_{r \rightarrow 0} \sup_{y \in B_r(\hat{x}) \setminus \{\hat{x}\}} (-\operatorname{div}(|D\phi|^{p-2}D\phi + a(y)|D\phi|^{p-2}D\phi)) \geq f(\hat{x}, u(\hat{x}), D\phi(\hat{x})).$$

Recalling that $u(\hat{x}) = u_\varepsilon(\hat{x})$, $Du_\varepsilon(\hat{x}) = 0 = D\phi(\hat{x})$, we obtain, from the above two displays,

$$0 \geq f(\hat{x}, u_\varepsilon(\hat{x}), Du_\varepsilon(\hat{x})) \geq f_\varepsilon(\hat{x}, u_\varepsilon(\hat{x}), Du_\varepsilon(\hat{x})),$$

as desired. □

Next, we show the convergence of inf-convolution u_ε in the Musielak-Orlicz-Sobolev space $W^{1,H(\cdot)}(\Omega)$ in the following two lemmas.

Lemma 3.5 *Under (1.3) and the preconditions of Lemma 3.1, we infer that the function $u \in W^{1,H(\cdot)}_{\text{loc}}(\Omega)$ and, up to a subsequence, $Du_\varepsilon \rightharpoonup Du$ weakly in $L^{H(\cdot)}(\Omega')$ for every $\Omega' \subset\subset \Omega$.*

Proof Take a cut-off function $\psi \in C_0^\infty(\Omega)$ fulfilling $\psi \equiv 1$ in Ω' and $0 \leq \psi \leq 1$ in Ω . Set ε so small that $\operatorname{supp} \psi =: E \subset \Omega_{r(\varepsilon)}$. Notice that $u_\varepsilon \rightarrow u$ locally uniformly in Ω , so we can define a test function

$$\varphi := (K - u_\varepsilon)\psi^q (\geq 0)$$

with $K := \sup_{\varepsilon; x \in \Omega'} |u_\varepsilon(x)|$ (finite). Hence, utilizing Lemma 3.3,

$$\begin{aligned}
 &\int_{\Omega_{r(\varepsilon)}} \varphi f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) dx + E(\varepsilon) \int_{\Omega_{r(\varepsilon)} \setminus \{Du_\varepsilon=0\}} \varphi dx \\
 &\leq \int_{\Omega_{r(\varepsilon)}} \langle |Du_\varepsilon|^{p-2}Du_\varepsilon + a(x)|Du_\varepsilon|^{q-2}Du_\varepsilon, D\varphi \rangle dx
 \end{aligned}$$

$$\begin{aligned}
 &= \int_{\Omega_{r(\varepsilon)}} q \psi^{q-1} (K - u_\varepsilon) (|Du_\varepsilon|^{p-2} Du_\varepsilon + a(x) |Du_\varepsilon|^{q-2} Du_\varepsilon, D\psi) dx \\
 &- \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, Du_\varepsilon) dx.
 \end{aligned}$$

Namely,

$$\begin{aligned}
 \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, Du_\varepsilon) dx &\leq q \int_{\Omega_{r(\varepsilon)}} \psi^{q-1} (K - u_\varepsilon) (|Du_\varepsilon|^{p-1} + a(x) |Du_\varepsilon|^{q-1}) |D\psi| dx \\
 &\quad + \int_{\Omega_{r(\varepsilon)}} \varphi |f_\varepsilon(x, u_\varepsilon, Du_\varepsilon)| dx + |E(\varepsilon)| \int_{\Omega_{r(\varepsilon)}} \varphi dx \\
 &=: I_1 + I_2 + I_3.
 \end{aligned}$$

Noting $q - 1 \geq \frac{q(p-1)}{p}$, $0 \leq \psi \leq 1$ and applying Young’s inequality with ϵ , we have

$$\begin{aligned}
 I_1 &\leq q \int_{\Omega_{r(\varepsilon)}} (K - u_\varepsilon) |D\psi| \psi^{\frac{q(p-1)}{p}} |Du_\varepsilon|^{p-1} + a(x) (K - u_\varepsilon) |D\psi| \psi^{q-1} |Du_\varepsilon|^{q-1} dx \\
 &\leq \epsilon \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, Du_\varepsilon) dx + C(q, \epsilon) \int_{\Omega_{r(\varepsilon)}} H(x, (K - u_\varepsilon) |D\psi|) dx.
 \end{aligned}$$

In view of the growth condition on f , we can deal with I_2 as

$$\begin{aligned}
 I_2 &\leq \int_{\Omega_{r(\varepsilon)}} (K - u_\varepsilon) \psi^q \gamma_\infty (|Du_\varepsilon|^{p-1} + a(x) |Du_\varepsilon|^{q-1}) + (K - u_\varepsilon) \psi^q \Phi dx \\
 &\leq \epsilon \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, Du_\varepsilon) dx + C(\gamma_\infty, \epsilon) \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, K - u_\varepsilon) dx \\
 &\quad + C(K, \|\Phi\|_{L^\infty(E)}, E).
 \end{aligned}$$

Here $\gamma_\infty := \max_{\tau \in [0, K]} \gamma(\tau)$. For I_3 , we have

$$I_3 \leq C(K, E),$$

where we assume $|E(\varepsilon)| \leq 1$ without loss of generality. Choosing proper $\epsilon \in (0, 1)$ and merging these above estimates yields that

$$\int_{\Omega'} H(x, Du_\varepsilon) dx \leq \int_{\Omega_{r(\varepsilon)}} \psi^q H(x, Du_\varepsilon) dx \leq C(p, q, a, K, \gamma, \Phi, D\psi, E).$$

This indicates that Du_ε is uniformly bounded in $L^{H(\cdot)}(\Omega')$ with respect to ε , which further deduces that there exists a function $Du \in L^{H(\cdot)}(\Omega')$ such that $Du_\varepsilon \rightharpoonup Du$ weakly in $L^{H(\cdot)}(\Omega')$ up to a subsequence owing to $L^{H(\cdot)}(\Omega')$ being a reflexible Banach space. Finally, u belongs to $W^{1, H(\cdot)}(\Omega')$ with Du as its weak derivative. \square

Lemma 3.6 *With (1.3) and the hypotheses of Lemma 3.1, we arrive at $u_\varepsilon \rightharpoonup u$ as $\varepsilon \rightarrow 0$, up to a subsequence, in $W^{1, H(\cdot)}(\Omega')$ for each $\Omega' \subset\subset \Omega$.*

Proof Take a cut-off function $\psi \in C_0^\infty(\Omega)$ satisfying $\psi \equiv 1$ in Ω' and $0 \leq \psi \leq 1$ in Ω . Let ε so small that $\text{supp } \psi =: E \subset \Omega_{r(\varepsilon)}$. Notice that $u_\varepsilon \leq u$ in Ω , so we define a test function

$$\varphi := (u - u_\varepsilon)\psi (\geq 0).$$

It is easy to find that $\varphi \in W_0^{1,H(\cdot)}(\Omega_{r(\varepsilon)})$. Employing again Lemma 3.3 obtains

$$\begin{aligned} & \int_{\Omega_{r(\varepsilon)}} \langle A(x, Du) - A(x, Du_\varepsilon), D\varphi \rangle dx \\ & \leq \int_{\Omega_{r(\varepsilon)}} \langle A(x, Du), D\varphi \rangle dx + \int_{\Omega_{r(\varepsilon)}} \varphi |f_\varepsilon(x, u_\varepsilon, Du_\varepsilon)| dx + |E(\varepsilon)| \int_{\Omega_{r(\varepsilon)}} \varphi dx. \end{aligned}$$

After manipulation, we get

$$\begin{aligned} & \int_{\Omega_{r(\varepsilon)}} \langle A(x, Du) - A(x, Du_\varepsilon), Du - Du_\varepsilon \rangle \psi dx \\ & \leq \int_E \psi \langle A(x, Du), Du - Du_\varepsilon \rangle dx + \int_E (u - u_\varepsilon) \langle A(x, Du_\varepsilon), D\psi \rangle dx \\ & \quad + \int_E \psi (u - u_\varepsilon) |f_\varepsilon(x, u_\varepsilon, Du_\varepsilon)| dx + |E(\varepsilon)| \int_E (u - u_\varepsilon) \psi dx \\ & = J_1 + J_2 + J_3 + J_4. \end{aligned}$$

For J_2 , via the inequalities (2.3) and (2.4), we derive

$$\begin{aligned} J_2 & \leq \|u - u_\varepsilon\|_{L^\infty(E)} \int_E |A(x, Du_\varepsilon)| |D\psi| dx \\ & \leq C \|u - u_\varepsilon\|_{L^\infty(E)} \int_E H^*(x, |A(x, Du_\varepsilon)|) + H(x, |D\psi|) dx \\ & \leq C \|u - u_\varepsilon\|_{L^\infty(E)} \int_E H(x, |Du_\varepsilon|) + H(x, |D\psi|) dx. \end{aligned}$$

Using the growth assumption on f and Young's inequality leads to

$$\begin{aligned} J_3 & \leq \|u - u_\varepsilon\|_{L^\infty(E)} \int_E \gamma_\infty (|Du_\varepsilon|^{p-1} + a(x)|Du_\varepsilon|^{q-1}) + \Phi dx \\ & \leq C \|u - u_\varepsilon\|_{L^\infty(E)} \left[\int_E H(x, |Du_\varepsilon|) dx + (1 + \|\Phi\|_{L^\infty(E)})|E| \right], \end{aligned}$$

where γ_∞ is the same as that in Lemma 3.5. As for J_4 ,

$$J_4 \leq |E(\varepsilon)| \|u - u_\varepsilon\|_{L^\infty(E)} |E|.$$

Recalling that $Du_\varepsilon \rightharpoonup Du$ weakly in $L^{H(\cdot)}(E)$ in Lemma 3.5 and $u_\varepsilon \rightarrow u$ locally uniformly in Proposition 2.5, we know $J_1 + J_2 + J_3 + J_4$ tends to 0. In other words,

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega'} \langle A(x, Du) - A(x, Du_\varepsilon), Du - Du_\varepsilon \rangle dx = 0.$$

Following the calculations in [23, page 9], we further arrive at

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega'} H(x, Du - Du_\varepsilon) dx = 0,$$

which implies the desired result. □

Finally, we end this section by verifying that the locally Lipschitz continuous viscosity supersolutions to (1.1) are also weak supersolutions, as stated in Theorem 1.1. We next intend to apply the previous convergence results to pass to the limit in the display of Lemma 3.3.

Proof of Theorem 1.1 Let $\varphi \in C_0^\infty(\Omega)$ be a nonnegative test function and set an open $\Omega' \subset\subset \Omega$ such that $\text{supp } \varphi \subset \Omega'$. Fix a sufficiently small $\varepsilon_0 > 0$ fulfilling $\Omega' \subset \Omega_{r(\varepsilon)}$ for $0 < \varepsilon < \varepsilon_0$. Now we want to show

$$\int_{\Omega} \langle |Du|^{p-2} Du + a(x)|Du|^{q-2} Du, D\varphi \rangle dx \geq \int_{\Omega} \varphi f(x, u, Du) dx, \tag{3.10}$$

which is the definition of weak supersolution of (1.1). This desired claim shall follow through Lemma 3.3, once the forthcoming displays are justified:

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \int_{\Omega'} \langle |Du_\varepsilon|^{p-2} Du_\varepsilon + a(x)|Du_\varepsilon|^{q-2} Du_\varepsilon, D\varphi \rangle dx \\ &= \int_{\Omega'} \langle |Du|^{p-2} Du + a(x)|Du|^{q-2} Du, D\varphi \rangle dx, \end{aligned} \tag{3.11}$$

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega'} \varphi f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) dx = \int_{\Omega'} \varphi f(x, u, Du) dx \tag{3.12}$$

and

$$\lim_{\varepsilon \rightarrow 0} E(\varepsilon) \int_{\Omega'} \varphi dx = 0. \tag{3.13}$$

First, the limit (3.13) is obviously valid. Next, we demonstrate the validity of (3.11). We shall employ the elementary vector inequality (see [38]):

$$||\xi_1|^{t-2}\xi_1 - |\xi_2|^{t-2}\xi_2| \leq \begin{cases} (t-1)|\xi_1 - \xi_2|(|\xi_1|^{t-2} + |\xi_2|^{t-2}), & \text{if } t \geq 2, \\ 2^{2-t}|\xi_1 - \xi_2|^{t-1}, & \text{if } 1 < t < 2, \end{cases} \tag{3.14}$$

where $\xi_1, \xi_2 \in \mathbb{R}^n$. We split the proof of (3.11) into three cases. □

For the case $1 < p \leq q < 2$, we use (2.3), (2.5)–(2.7), the basic inequality (3.14) to get

$$\begin{aligned} & \int_{\Omega'} (|Du_\varepsilon|^{p-2} Du_\varepsilon - |Du|^{p-2} Du + a(x)(|Du_\varepsilon|^{q-2} Du_\varepsilon - |Du|^{q-2} Du), D\varphi) dx \\ & \leq \int_{\Omega'} (|Du_\varepsilon|^{p-2} Du_\varepsilon - |Du|^{p-2} Du + a(x)(|Du_\varepsilon|^{q-2} Du_\varepsilon - |Du|^{q-2} Du)) |D\varphi| dx \\ & \leq C \int_{\Omega'} (|Du_\varepsilon - Du|^{p-1} + a(x)|Du_\varepsilon - Du|^{q-1}) |D\varphi| dx \\ & \leq C \|h(x, Du_\varepsilon - Du)\|_{L^{H^*(\cdot)}(\Omega')} \|D\varphi\|_{L^{H(\cdot)}(\Omega')} \\ & \leq C \|D\varphi\|_{L^{H(\cdot)}(\Omega')} \max \left\{ (\varrho_{H^*}(h(x, Du_\varepsilon - Du)))^{\frac{q}{q+1}}, (\varrho_{H^*}(h(x, Du_\varepsilon - Du)))^{\frac{p}{p+1}} \right\} \\ & \leq C \|D\varphi\|_{L^{H(\cdot)}(\Omega')} \max \left\{ \left(\int_{\Omega'} H(x, Du_\varepsilon - Du) dx \right)^{\frac{q}{q+1}}, \left(\int_{\Omega'} H(x, Du_\varepsilon - Du) dx \right)^{\frac{p}{p+1}} \right\} \\ & \rightarrow 0 \end{aligned}$$

as $\varepsilon \rightarrow 0$, by applying $Du_\varepsilon \rightarrow Du$ in $L^{H(\cdot)}(\Omega')$ in Lemma 3.6. Here $h(x, z) := |z|^{p-1} + a(x)|z|^{q-1}$.

When $2 \leq p \leq q$, exploiting again (2.3), (2.5)–(2.7) as well as (3.14), and applying $Du_\varepsilon \rightarrow Du$ in $L^{H(\cdot)}(\Omega')$, we can see that

$$\begin{aligned} & \int_{\Omega'} (|Du_\varepsilon|^{p-2} Du_\varepsilon - |Du|^{p-2} Du + a(x)(|Du_\varepsilon|^{q-2} Du_\varepsilon - |Du|^{q-2} Du), D\varphi) dx \\ & \leq C \int_{\Omega'} [(|Du_\varepsilon|^{p-2} + |Du|^{p-2}) + a(x)(|Du_\varepsilon|^{q-2} + |Du|^{q-2})] |Du_\varepsilon - Du| |D\varphi| dx \\ & \leq C \|D\varphi\|_{L^\infty(\Omega')} \int_{\Omega'} [(1 + |Du_\varepsilon|^{p-1} + |Du|^{p-1}) + a(x)(1 + |Du_\varepsilon|^{q-1} + |Du|^{q-1})] \\ & \quad \times |Du_\varepsilon - Du| dx \\ & = C \|D\varphi\|_{L^\infty(\Omega')} \int_{\Omega'} [(1 + a(x)) + h(x, Du_\varepsilon) + h(x, Du)] |Du_\varepsilon - Du| dx \\ & \leq C \|Du_\varepsilon - Du\|_{L^{H(\cdot)}(\Omega')} + C \|h(x, Du)\|_{L^{H^*(\cdot)}(\Omega')} \|Du_\varepsilon - Du\|_{L^{H(\cdot)}(\Omega')} \\ & \quad + C \|h(x, Du_\varepsilon)\|_{L^{H^*(\cdot)}(\Omega')} \|Du_\varepsilon - Du\|_{L^{H(\cdot)}(\Omega')} \\ & \rightarrow 0 \end{aligned}$$

as $\varepsilon \rightarrow 0$, where the constant C depends on $p, q, \|a\|_{L^\infty(\Omega')}, \|D\varphi\|_{L^\infty(\Omega')}, |\Omega|$. Here we need to notice that the quantity $\int_{\Omega'} H(x, Du_\varepsilon) dx$ is uniformly bounded.

In the last case $1 < p < 2 \leq q$, we combine the previous two scenarios to deduce the claim (3.11). Specifically,

$$\begin{aligned} & \int_{\Omega'} (|Du_\varepsilon|^{p-2} Du_\varepsilon - |Du|^{p-2} Du + a(x)(|Du_\varepsilon|^{q-2} Du_\varepsilon - |Du|^{q-2} Du), D\varphi) dx \\ & \leq C \int_{\Omega'} [|Du_\varepsilon - Du|^{p-1} + a(x)(|Du_\varepsilon|^{q-2} + |Du|^{q-2}) |Du_\varepsilon - Du|] |D\varphi| dx \\ & \leq C \int_{\Omega'} (|Du_\varepsilon - Du|^{p-1} + a(x)|Du_\varepsilon - Du|^{q-1}) |D\varphi| dx \end{aligned}$$

$$\begin{aligned}
 &+ C \int_{\Omega'} a(x)(1 + |Du_\varepsilon|^{q-1} + |Du|^{q-1})|Du_\varepsilon - Du|D\varphi| dx \\
 \leq &C \int_{\Omega'} h(x, Du_\varepsilon - Du)|D\varphi| dx + C \int_{\Omega'} |Du_\varepsilon - Du| dx \\
 &+ C \int_{\Omega'} h(x, Du_\varepsilon)|Du_\varepsilon - Du| dx + C \int_{\Omega'} h(x, Du)|Du_\varepsilon - Du| dx,
 \end{aligned}$$

which tends to 0 as $\varepsilon \rightarrow 0$, where C depends upon $p, q, \|a\|_{L^\infty(\Omega')}, \|D\varphi\|_{L^\infty(\Omega')}, |\Omega'|$.

Eventually, let us prove the claim (3.12). Via the uniform continuity of f , for each $\epsilon > 0$, there is a $\delta > 0$ depending only on ϵ such that

$$|f(x, u_\varepsilon(x), Du_\varepsilon(x)) - f(y, u_\varepsilon(x), Du_\varepsilon(x))| \leq \epsilon \quad \text{for } y \in B_\delta(x).$$

Now select $\epsilon'_0 > 0$ to satisfy $r(\epsilon) < \delta$ if $0 < \epsilon < \epsilon'_0$. Then we have

$$f(x, u_\varepsilon(x), Du_\varepsilon(x)) < \epsilon + f(y, u_\varepsilon(x), Du_\varepsilon(x))$$

for each $x \in \Omega'$ and $y \in B_{r(\epsilon)}(x)$. In particular,

$$f(x, u_\varepsilon(x), Du_\varepsilon(x)) < \epsilon + f_\varepsilon(x, u_\varepsilon(x), Du_\varepsilon(x))$$

by the definition of f_ε , and moreover

$$0 \leq |f(x, u_\varepsilon(x), Du_\varepsilon(x)) - f_\varepsilon(x, u_\varepsilon(x), Du_\varepsilon(x))| \leq \epsilon.$$

We thus have

$$\int_{\Omega'} |f(x, u_\varepsilon, Du_\varepsilon) - f_\varepsilon(x, u_\varepsilon, Du_\varepsilon)|\varphi dx \leq \epsilon \|\varphi\|_{L^\infty(\Omega')} |\Omega'|.$$

In view of (1) in Proposition 2.5, it is known that $\|u_\varepsilon\|_{L^\infty(\Omega')} \leq \|u\|_{L^\infty(\Omega')}$ for any ε . Namely,

$$\max_{\tau \in [0, \|u_\varepsilon\|_{L^\infty(\Omega')}] } \{\gamma(\tau)\} \leq \max_{\tau \in [0, \|u\|_{L^\infty(\Omega')}] } \{\gamma(\tau)\} =: \gamma_\infty.$$

According to the growth condition on f ,

$$|f(x, u_\varepsilon, Du)| \leq \gamma_\infty(|Du|^{p-1} + a(x)|Du|^{q-1}) + \Phi(x) \quad \text{in } \Omega',$$

which belongs to $L^{H^*(\cdot)}(\Omega')$. Thereby we can exploit the dominated convergence theorem to infer

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega'} |f(x, u_\varepsilon, Du) - f(x, u, Du)|\varphi dx = 0.$$

We finally address the term $\int_{\Omega'} |f(x, u_\varepsilon, Du_\varepsilon) - f(x, u_\varepsilon, Du)|\varphi dx$,

$$\begin{aligned} & \int_{\Omega'} |f(x, u_\varepsilon, Du_\varepsilon) - f(x, u_\varepsilon, Du)|\varphi dx \\ & \leq C \int_{\Omega'} |Du_\varepsilon - Du|\varphi dx \\ & \leq C \|\varphi\|_{L^{H^*(\cdot)}(\Omega')} \|Du_\varepsilon - Du\|_{L^{H(\cdot)}(\Omega')} \\ & \rightarrow 0 \text{ as } \varepsilon \rightarrow 0, \end{aligned}$$

where we have utilized the Lipschitz continuity of f in the third variable, and the convergence $Du_\varepsilon \rightarrow Du$ in $L^{H(\cdot)}(\Omega')$.

Merging these above estimates yields that

$$\begin{aligned} & \int_{\Omega'} |f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) - f(x, u, Du)|\varphi dx \\ & \leq \int_{\Omega'} |f_\varepsilon(x, u_\varepsilon, Du_\varepsilon) - f(x, u_\varepsilon, Du_\varepsilon)|\varphi dx \\ & \quad + \int_{\Omega'} |f(x, u_\varepsilon, Du_\varepsilon) - f(x, u_\varepsilon, Du)|\varphi dx \\ & \quad + \int_{\Omega'} |f(x, u_\varepsilon, Du) - f(x, u, Du)|\varphi dx \end{aligned}$$

converges to 0 by sending $\varepsilon \rightarrow 0$. Hereto, we have verified the displays (3.11)–(3.13), from which we get the inequality (3.10).

4 weak solutions are viscosity solutions

In this section, we prove that weak supersolutions are viscosity supersolutions to (1.1), that is Theorem 1.4, by using comparison principle for weak solutions, subsequently giving three examples of comparison results. Then weak subsolutions can be showed to be viscosity subsolutions in a similar way.

Proof of Theorem 1.4 We argue by contradiction. If not, there exists a $\varphi \in C^2(\Omega)$ touching u from below at $x_0 \in \Omega$, that is,

$$\begin{cases} u(x_0) = \varphi(x_0), \\ u(x) > \varphi(x) \text{ for } x \neq x_0, \\ D\varphi(x_0) \neq 0, \end{cases}$$

and however,

$$-\operatorname{div} A(x_0, D\varphi(x_0)) < f(x_0, u(x_0), D\varphi(x_0)).$$

By means of continuity, for some $\delta > 0$ there is a small enough $r > 0$ such that

$$-\operatorname{div} A(x, D\varphi(x)) \leq f(x, \varphi, D\varphi) - \delta \tag{4.1}$$

for $x \in B_r := B_r(x_0)$. Denote $\tilde{\varphi} := \varphi + m$ with $m > 0$ to be chosen later. In view of (4.1), we have

$$\begin{aligned} & -\operatorname{div} A(x, D\tilde{\varphi}) - f(x, \tilde{\varphi}, D\tilde{\varphi}) \\ &= -\operatorname{div} A(x, D\varphi) - f(x, \varphi, D\varphi) + f(x, \varphi, D\varphi) - f(x, \tilde{\varphi}, D\tilde{\varphi}) \\ &\leq -\delta + f(x, \varphi, D\varphi) - f(x, \varphi + m, D\varphi) \\ &\leq -\frac{\delta}{2}, \end{aligned}$$

if $m > 0$ is sufficiently small. Indeed, we can pick $m \in (0, \frac{1}{2} \min_{\partial B_r} \{u - \varphi\})$ (note the lower semicontinuity of u) small, taking into account the uniform continuity of f , such that

$$|f(x, \varphi, D\varphi) - f(x, \varphi + m, D\varphi)| \leq \frac{\delta}{2}.$$

Hence, $\tilde{\varphi}$ is a weak subsolution to (1.1) in B_r , as well. Observe that $\tilde{\varphi} = \varphi + m < \varphi + u - \varphi = u$ on ∂B_r . Thereby through the (CPP) we get $u \geq \tilde{\varphi}$ in B_r . Nonetheless, $u(x_0) = \varphi(x_0) < \varphi(x_0) + m$, which is a contradiction. Then u is a viscosity supersolution to (1.1). □

A fundamental issue in Theorem 1.4 is the availability of comparison principle for weak solutions. Nevertheless, it is not undemanding to establish such principle for the nonhomogeneous double phase equations with very general structure. Hence, we next for three slightly special cases prove comparison principle for weak solutions.

Lemma 4.1 *Assume that $u, v \in W^{1,H(\cdot)}(\Omega)$ are the weak subsolution and supersolution, respectively, to $-\operatorname{div} A(x, Dw) = f(x, w)$ in Ω , where f is decreasing in the w -variable. If $u \leq v$ on $\partial\Omega$, then $u \leq v$ a.e. in Ω .*

Proof Since u, v are separately weak subsolution and supersolution such that $u \leq v$ on $\partial\Omega$, the auxiliary function

$$w := (u - v - l)_+, \quad l > 0,$$

can be chosen as a test function, which belongs to the space $W_0^{1,H(\cdot)}(\Omega)$. Therefore, we get

$$\int_{\Omega} \langle A(x, Du) - A(x, Dv), Dw \rangle dx \leq \int_{\Omega} (f(x, u) - f(x, v))w dx.$$

We use the fact that f is decreasing with respect to the second variable to arrive at

$$\int_{\Omega} (f(x, u) - f(x, v))(u - v - l)_+ dx \leq 0.$$

Furthermore, due to the strictly monotone increasing property of the operator $A(x, \cdot)$, i.e., $\langle A(x, \xi) - A(x, \zeta), \xi - \zeta \rangle > 0$ for $\xi \neq \zeta \in \mathbb{R}^n$, it follows that

$$0 \leq \int_{\Omega \cap \{(u-v-l)_+ > 0\}} \langle A(x, Du) - A(x, Dv), Du - Dv \rangle dx \leq 0.$$

This implies that $w \equiv 0$ a.e. in Ω . In other words, $u \leq v + l$ a.e. in Ω . Letting $l \rightarrow 0$, we can see that $u \leq v$ a.e. in Ω . □

Lemma 4.2 *Suppose that $f(x, \tau, \eta)$ is decreasing in τ , and is locally Lipschitz continuous with respect to η in $\Omega \times \mathbb{R} \times \mathbb{R}^n$. Let u, v be the weak subsolution and supersolution respectively to (1.1) such that*

$$|Du(x)| + |Dv(x)| \geq \delta \quad \text{a.e. } x \in \Omega$$

with $\delta > 0$ any number. Let also $2 \leq p \leq q$. Then there is $\varepsilon > 0$ such that, for every domain $E \subset \subset \Omega$ fulfilling $|E| \leq \varepsilon$, whenever $u \leq v$ on ∂E then it holds that $u \leq v$ a.e. in E .

Proof Selecting $w = (u - v)_+ \chi_E \in W_0^{1,H(\cdot)}(\Omega)$ as a test function in the weak formulation of (1.1), we derive, from the hypotheses on f ,

$$\begin{aligned} & \int_{\Omega} \langle A(x, Du) - A(x, Dv), Dw \rangle dx \\ & \leq \int_{\Omega} (f(x, u, Du) - f(x, v, Dv))w dx \\ & = \int_{\Omega} (f(x, u, Du) - f(x, v, Du))(u - v)_+ \chi_E + (f(x, v, Du) \\ & \quad - f(x, v, Dv))(u - v)_+ \chi_E dx \\ & \leq \int_{\Omega} (f(x, v, Du) - f(x, v, Dv))(u - v)_+ \chi_E dx \\ & \leq C \int_E |Du - Dv|(u - v)_+ dx. \end{aligned} \tag{4.2}$$

Now by the basic inequality (see [38, page 97])

$$C(|\xi| + |\eta|)^{p-2} |\xi - \eta|^2 \leq (|\xi|^{p-2} \xi - |\eta|^{p-2} \eta) \cdot (\xi - \eta) \quad \text{for } p \geq 2,$$

(4.2) turns into

$$\int_E (|Du| + |Dv|)^{p-2} |D(u - v)_+|^2 + a(x)(|Du| + |Dv|)^{q-2} |D(u - v)_+|^2 dx$$

$$\begin{aligned}
 &\leq C \int_E |Du - Dv|(u - v)_+ dx \\
 &\leq C(E) \int_E |D(u - v)_+|^2 dx \\
 &\leq C(E) \int_E (|Du| + |Dv|)^{2-p} (|Du| + |Dv|)^{p-2} |D(u - v)_+|^2 dx \\
 &\leq \delta^{2-p} C(E) \int_E (|Du| + |Dv|)^{p-2} |D(u - v)_+|^2 dx,
 \end{aligned}$$

where in the third line we utilized the Hölder and Poincaré inequalities, and the quantity $C(E)$ will tend to 0 when the measure $|E|$ goes to 0. From above, we can find that if the domain E is small enough, the last inequality is self-contradictory. That is, $(u - v)_+ \equiv 0$ in E , which leads to $u \leq v$ a.e. in E . □

Remark 4.3 Lemma 4.1 could be exploited in the proof of Theorem 1.4 apparently. Observe that $D\varphi(x_0) \neq 0$ and $\varphi \in C^2(\Omega)$, so we can take such small ball $B_r(x_0)$ that $|B_r(x_0)| \leq \varepsilon$ (ε is provided by Lemma 4.2) and moreover $|D\varphi(x)| \geq \delta > 0$ in $B_r(x_0)$. Thus we can keep track of the proof of Theorem 1.4 to deduce that a lower semicontinuous weak supersolution is a viscosity supersolution to (1.1) by Lemma 4.2.

Lemma 4.4 *Suppose that $f(x, \tau, \eta)$ is decreasing in τ , and is locally Lipschitz continuous with respect to η in $\Omega \times \mathbb{R} \times \mathbb{R}^n$. Let $1 < p < 2$ and $u, v \in W_{loc}^{1,\infty}(\Omega)$ be the weak subsolution and supersolution respectively to (1.1). Then there is $\varepsilon > 0$ such that, for every domain $E \subset\subset \Omega$ fulfilling $|E| \leq \varepsilon$, whenever $u \leq v$ on ∂E then it holds that $u \leq v$ a.e. in E .*

Proof Selecting $w = (u - v)_+ \chi_E \in W_0^{1,H^{(\cdot)}}(\Omega)$ as a test function, we can also derive the estimate (4.2). Now by using the basic inequality (see [38, page 100])

$$C(1 + |\xi|^2 + |\eta|^2)^{\frac{p-2}{2}} |\xi - \eta|^2 \leq (|\xi|^{p-2}\xi - |\eta|^{p-2}\eta) \cdot (\xi - \eta) \quad \text{for } 1 < p < 2,$$

and

$$0 \leq a(x)(|\xi|^{q-2}\xi - |\eta|^{q-2}\eta) \cdot (\xi - \eta) \quad \text{for } q > 1,$$

(4.2) becomes

$$\begin{aligned}
 &\int_E (1 + |Du|^2 + |Dv|^2)^{\frac{p-2}{2}} |D(u - v)_+|^2 dx \\
 &\leq C \int_E |Du - Dv|(u - v)_+ dx \\
 &\leq C(E) \int_E |D(u - v)_+|^2 dx \\
 &\leq C(E) \int_E (1 + |Du|^2 + |Dv|^2)^{\frac{2-p}{2}} (1 + |Du|^2 + |Dv|^2)^{\frac{p-2}{2}} |D(u - v)_+|^2 dx
 \end{aligned}$$

$$\leq (1 + \|Du\|_{L^\infty(E)} + \|Dv\|_{L^\infty(E)})^{2-p} C(E) \int_E (1 + |Du|^2 + |Dv|^2)^{\frac{p-2}{2}} |D(u-v)_+|^2 dx,$$

where the quantity $C(E)$ will tend to 0 when the measure $|E|$ goes to 0. From the estimate above, we find that if the domain E is small enough, the last inequality is self-contradictory. That is, $(u - v)_+ \equiv 0$ in E , which leads to $u \leq v$ a.e. in E . \square

5 Lipschitz continuity of viscosity solutions

We in this part show that the bounded viscosity solutions to (1.1) are locally Lipschitz continuous. The strategy is to verify first the Hölder continuity of viscosity solutions by using the Ishii–Lions methods, and further, based on the Hölder continuity, to demonstrate the Lipschitz continuity of viscosity solutions through the Ishii-Lions methods again. The similar idea can be found for instance in [2]. For the sake of convenience, we suppose the domain Ω is a unit ball B_1 .

Lemma 5.1 (Local Hölder continuity) *Let u be a bounded viscosity solution to (1.1) in B_1 . Assume that $0 \leq a(x) \in C^1(B_1)$, $p \leq q \leq p + 1$ and (1.3) are in force. Then for each $\beta \in (0, 1)$, there exists a constant $C > 0$, depending on $n, p, q, \beta, \gamma_\infty, \|a\|_{C^1(B_1)}, \|u\|_{L^\infty(B_1)}$ and $\|\Phi\|_{L^\infty(B_1)}$, such that*

$$|u(x) - u(y)| \leq C|x - y|^\beta$$

for any $x, y \in B_{3/4}$, where $\gamma_\infty := \max_{t \in [0, \|u\|_{L^\infty(B_1)}]} \gamma(t)$.

Proof Fix $x_0, y_0 \in B_{3/4}$. We now aim at showing that there are two proper constants $L_1, L_2 > 0$ such that

$$\omega := \sup_{x, y \in B_{3/4}} \left(u(x) - u(y) - L_1\phi(|x - y|) - \frac{L_2}{2}|x - x_0|^2 - \frac{L_2}{2}|y - y_0|^2 \right) \leq 0, \tag{5.1}$$

where $\phi(r) = r^\beta$ with $\beta \in (0, 1)$.

To this end, assume on the contrary that (5.1) is not true and moreover $(\bar{x}, \bar{y}) \in \overline{B_{3/4} \times B_{3/4}}$ stands for the point where the supremum is achieved. We can easily know two facts that $\bar{x} \neq \bar{y}$ by $\omega > 0$, and $\bar{x}, \bar{y} \in B_{3/4}$ by choosing

$$L_2 \geq \frac{64\|u\|_{L^\infty(B_1)}}{(\min\{\text{dist}(x_0, \partial B_{3/4}), \text{dist}(y_0, \partial B_{3/4})\})^2}.$$

Besides,

$$|\bar{x} - \bar{y}| \leq \left(\frac{2\|u\|_{L^\infty(B_1)}}{L_1} \right)^{\frac{1}{\beta}}$$

is small enough, provided that L_1 is sufficiently large, which shall be utilized later.

Now by invoking the maximum principle for semicontinuous functions [16, Theorem 3.2], there are $X, Y \in \mathcal{S}^n$ such that

$$(\eta_1, X + L_2I) \in \bar{J}^{2,+}u(\bar{x}) \quad \text{and} \quad (\eta_2, Y - L_2I) \in \bar{J}^{2,-}u(\bar{y})$$

with

$$\begin{aligned} \eta_1 &= L_1D_x\phi(|\bar{x} - \bar{y}|) + L_2(\bar{x} - x_0) = L_1\phi'(|\bar{x} - \bar{y}|)\frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} + L_2(\bar{x} - x_0), \\ \eta_2 &= -L_1D_y\phi(|\bar{x} - \bar{y}|) - L_2(\bar{y} - y_0) = L_1\phi'(|\bar{x} - \bar{y}|)\frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} - L_2(\bar{y} - y_0). \end{aligned}$$

Via selecting $L_1 \geq C(\beta)L_2$ large enough, there holds that

$$\frac{\beta L_1}{2}|\bar{x} - \bar{y}|^{\beta-1} \leq |\eta_1|, |\eta_2| \leq 2\beta L_1|\bar{x} - \bar{y}|^{\beta-1}. \tag{5.2}$$

Furthermore, applying [15, Theorem 12.2], for any $\tau > 0$ such that $\tau Z < I$, we obtain

$$-\frac{2}{\tau} \begin{pmatrix} I & \\ & I \end{pmatrix} \leq \begin{pmatrix} X & \\ & -Y \end{pmatrix} \leq \begin{pmatrix} Z^\tau & -Z^\tau \\ -Z^\tau & Z^\tau \end{pmatrix}, \tag{5.3}$$

where

$$\begin{aligned} Z &= L_1\phi''(|\bar{x} - \bar{y}|)\frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \otimes \frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} + \frac{L_1\phi'(|\bar{x} - \bar{y}|)}{|\bar{x} - \bar{y}|} \left(I - \frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \otimes \frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \right) \\ &= \beta L_1|\bar{x} - \bar{y}|^{\beta-2} \left(I + (\beta - 2)\frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \otimes \frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \right) \end{aligned}$$

and

$$Z^\tau = (I - \tau Z)^{-1}Z$$

with $(I - \tau Z)^{-1}$ denoting the inverse of the matrix $I - \tau Z$. Now pick $\tau = \frac{1}{2\beta L_1|\bar{x} - \bar{y}|^{\beta-2}}$ such that

$$Z^\tau = 2\beta L_1|\bar{x} - \bar{y}|^{\beta-2} \left(I - 2\frac{\beta - 1}{3 - \beta}\frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \otimes \frac{\bar{x} - \bar{y}}{|\bar{x} - \bar{y}|} \right).$$

Observe that

$$\langle Z^\tau \xi, \xi \rangle = 2\beta \frac{\beta - 1}{3 - \beta} L_1|\bar{x} - \bar{y}|^{\beta-2} < 0 \tag{5.4}$$

for $\xi = \frac{\bar{x}-\bar{y}}{|\bar{x}-\bar{y}|}$. In addition, it follows from (5.3) that $X \leq Y$ and

$$\|X\|, \|Y\| \leq 4\beta L_1 |\bar{x} - \bar{y}|^{\beta-2}. \tag{5.5}$$

Let

$$A_s(\eta) := I + (s - 2) \frac{\eta}{|\eta|} \otimes \frac{\eta}{|\eta|} \text{ for } \eta \in \mathbb{R}^n \setminus \{0\}$$

with $s \in \{p, q\}$. An obvious fact is that the eigenvalues of $A_s(\eta)$ belong to the interval $[\min\{1, s - 1\}, \max\{1, s - 1\}]$. Now since u is a viscosity solution to (1.1), we have

$$F(\bar{x}, \eta_1, X + L_2I) - f(\bar{x}, u(\bar{x}), \eta_1) \leq 0$$

and

$$F(\bar{y}, \eta_2, Y - L_2I) - f(\bar{y}, u(\bar{y}), \eta_2) \geq 0.$$

Adding these two inequalities becomes

$$\begin{aligned} 0 &\leq |\eta_1|^{p-2} \text{tr}(A_p(\eta_1)(X + L_2I)) - |\eta_2|^{p-2} \text{tr}(A_p(\eta_2)(Y - L_2I)) \\ &\quad + a(\bar{x})|\eta_1|^{q-2} \text{tr}(A_q(\eta_1)(X + L_2I)) - a(\bar{y})|\eta_2|^{q-2} \text{tr}(A_q(\eta_2)(Y - L_2I)) \\ &\quad + |\eta_1|^{q-2} \eta_1 \cdot Da(\bar{x}) - |\eta_2|^{q-2} \eta_2 \cdot Da(\bar{y}) + f(\bar{x}, u(\bar{x}), \eta_1) - f(\bar{y}, u(\bar{y}), \eta_2) \\ &=: I_1 + I_2 + I_3 + I_4. \end{aligned} \tag{5.6}$$

In what follows, our goal is to justify $I_1 + I_2 + I_3 + I_4 < 0$ under suitable conditions, which reaches a contradiction so that the claim (5.1) is precisely true. First, we examine the term I_2 as

$$\begin{aligned} I_2 &= (a(\bar{x}) - a(\bar{y}))|\eta_1|^{q-2} \text{tr}(A_q(\eta_1)(X + L_2I)) \\ &\quad + a(\bar{y})(|\eta_1|^{q-2} - |\eta_2|^{q-2}) \text{tr}(A_q(\eta_1)(X + L_2I)) \\ &\quad + a(\bar{y})|\eta_2|^{q-2} [\text{tr}(A_q(\eta_1)(X + L_2I)) - \text{tr}(A_q(\eta_2)(Y - L_2I))] \\ &= (a(\bar{x}) - a(\bar{y}))|\eta_1|^{q-2} \text{tr}(A_q(\eta_1)(X + L_2I)) \\ &\quad + a(\bar{y})(|\eta_1|^{q-2} - |\eta_2|^{q-2}) \text{tr}(A_q(\eta_1)(X + L_2I)) \\ &\quad + a(\bar{y})|\eta_2|^{q-2} \text{tr}(A_q(\eta_1)(X - Y)) + a(\bar{y})|\eta_2|^{q-2} \text{tr}((A_q(\eta_1) - A_q(\eta_2))Y) \\ &\quad + L_2 a(\bar{y})|\eta_2|^{q-2} (\text{tr}(A_q(\eta_1)) + \text{tr}(A_q(\eta_2))) \\ &=: I_{21} + I_{22} + I_{23} + I_{24} + I_{25}. \end{aligned}$$

Let us mention that the treatment of I_2 is similar to that of J_1 in [24, Lemma 6.1]. However, we here give some details for the sake of readability. For I_{21} , by virtue of (5.2), (5.5) and the eigenvalues of $A_q(\eta_1)$, we have

$$I_{21} \leq C(\beta, q, \|a\|_{C^1(B_1)}) |\bar{x} - \bar{y}| L_1^{q-2} |\bar{x} - \bar{y}|^{(\beta-1)(q-2)} (n \|A_q(\eta_1)\| \|X\| + n L_2 \|A_q(\eta_1)\|)$$

$$\leq C(n, \beta, q, \|a\|_{C^1(B_1)})L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)+1}(L_1|\bar{x} - \bar{y}|^{\beta-2} + L_2).$$

We now consider I_{22} ,

$$\begin{aligned} I_{22} &\leq C(q, \beta)a(\bar{y})L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}L_2(n\|A_q(\eta_1)\| \|X\| + nL_2\|A_q(\eta_1)\|) \\ &\leq C(n, \beta, q)a(\bar{y})L_2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}(L_1|\bar{x} - \bar{y}|^{\beta-2} + L_2). \end{aligned}$$

Indeed, we evaluate $|\eta_1|^{q-2} - |\eta_2|^{q-2}$ as

$$\begin{aligned} |\eta_1|^{q-2} - |\eta_2|^{q-2} &= (q - 2)|\xi_{12}|^{q-3}(|\eta_1| - |\eta_2|) \\ &\leq |q - 2|C(q, \beta)L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}|\eta_1 - \eta_2| \\ &\leq C(q, \beta)L_2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}, \end{aligned}$$

where $|\xi_{12}|$ is between $|\eta_1|$ and $|\eta_2|$, and we have used (5.2) and $|\eta_1 - \eta_2| \leq 4L_2$. Due to (5.3) and (5.4), any eigenvalue of $X - Y$ is non-positive and at least one eigenvalue, denoted by $\bar{\lambda}(X - Y)$, is less than or equal to $8\beta\frac{\beta-1}{3-\beta}L_1|\bar{x} - \bar{y}|^{\beta-2}$. Next, we deal with the term I_{23} as

$$\begin{aligned} I_{23} &\leq a(\bar{y})|\eta_2|^{q-2} \sum_{i=1}^n \lambda_i(A_q(\eta_1))\lambda_i(X - Y) \\ &\leq a(\bar{y})|\eta_2|^{q-2} \min\{1, q - 1\}\bar{\lambda}(X - Y) \\ &\leq a(\bar{y})C(q, \beta)L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)} \left(-8\beta\frac{1-\beta}{3-\beta}\right) L_1|\bar{x} - \bar{y}|^{\beta-2} \\ &\leq -C(q, \beta)a(\bar{y})L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} \end{aligned}$$

with $\lambda_1(A_q(\eta_1)) \leq \lambda_2(A_q(\eta_1)) \leq \dots \leq \lambda_n(A_q(\eta_1))$ and $\lambda_1(X - Y) \leq \lambda_2(X - Y) \leq \dots \leq \lambda_n(X - Y)$. Here we need to note $0 < \beta < 1$. In order to evaluate I_{24} , we first notice

$$\begin{aligned} \|A_q(\eta_1) - A_q(\eta_2)\| &\leq 2|q - 2| \left| \frac{\eta_1}{|\eta_1|} - \frac{\eta_2}{|\eta_2|} \right| \\ &\leq 2|q - 2| \max \left\{ \frac{|\eta_1 - \eta_2|}{|\eta_1|}, \frac{|\eta_1 - \eta_2|}{|\eta_2|} \right\} \\ &\leq C(q, \beta) \frac{L_2}{L_1|\bar{x} - \bar{y}|^{\beta-1}}. \end{aligned}$$

Thereby, from (5.2), (5.5) and the preceding inequality,

$$\begin{aligned} I_{24} &\leq a(\bar{y})|\eta_2|^{q-2}n\|A_q(\eta_1) - A_q(\eta_2)\| \|Y\| \\ &\leq C(n, q, \beta)a(\bar{y})(L_1|\bar{x} - \bar{y}|^{\beta-1})^{q-2} \frac{L_2}{L_1|\bar{x} - \bar{y}|^{\beta-1}} L_1|\bar{x} - \bar{y}|^{\beta-2} \\ &= C(n, q, \beta)a(\bar{y})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)-1}. \end{aligned}$$

The term I_{25} finally can be treated by

$$I_{25} \leq a(\bar{y})|\eta_2|^{q-2}L_22n \max\{1, q - 1\} \leq C(n, q, \beta)a(\bar{y})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)}.$$

Now merging these estimates of I_{21} – I_{25} yields that

$$\begin{aligned} I_2 \leq & -C(q, \beta)a(\bar{y})L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} \\ & + C(n, q, \beta)a(\bar{y})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)-1} \\ & + C(n, q, \beta)a(\bar{y})L_2^2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)} \\ & + C(n, q, \beta, \|a\|_{C^1(B_1)})L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)} \\ & + C(n, q, \beta, \|a\|_{C^1(B_1)})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)+1}, \end{aligned}$$

where we have employed the fact that $|\bar{x} - \bar{y}| < 1$. Analogously, we could derive

$$\begin{aligned} I_1 \leq & -C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} + C(n, p, \beta)L_2L_1^{p-2}|\bar{x} - \bar{y}|^{(\beta-1)(p-2)-1} \\ & + C(n, p, \beta)L_2^2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(p-3)}, \end{aligned}$$

where we just note $a(\cdot) \equiv 1$. The term I_3 is directly estimated as

$$I_3 \leq |\eta_1|^{q-1}|Da(\bar{x})| + |\eta_2|^{q-1}|Da(\bar{y})| \leq C(q, \beta, \|a\|_{C^1(B_1)})L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)}$$

by applying (5.2). As for I_4 , according to the growth condition on f ,

$$\begin{aligned} I_4 \leq & |f(\bar{x}, u(\bar{x}), \eta_1)| + |f(\bar{y}, u(\bar{y}), \eta_2)| \\ \leq & \gamma(|u(\bar{x})|)(|\eta_1|^{p-1} + a(\bar{x})|\eta_1|^{q-1}) + \Phi(\bar{x}) \\ & + \gamma(|u(\bar{y})|)(|\eta_2|^{p-1} + a(\bar{y})|\eta_2|^{q-1}) + \Phi(\bar{y}) \\ \leq & \gamma_\infty C(p, q, \beta)(L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)} \\ & + \|a\|_{L^\infty(B_1)}L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)}) \\ & + \|\Phi\|_{L^\infty(B_1)} \end{aligned}$$

with $\gamma_\infty := \max_{t \in [0, \|u\|_{L^\infty(B_1)}]} \gamma(t)$.

We eventually gather the estimates on I_1 – I_4 with (5.6) to infer that

$$\begin{aligned} 0 \leq & [-C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} + C(n, p, \beta)L_2L_1^{p-2}|\bar{x} - \bar{y}|^{(\beta-1)(p-2)-1} \\ & + C(n, p, \beta)L_2^2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(p-3)} + C(p, q, \beta, \gamma_\infty)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)} \\ & + C(n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty)L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)} + \|\Phi\|_{L^\infty(B_1)} \\ & + C(n, q, \beta, \|a\|_{C^1(B_1)})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)+1}] \\ & + a(\bar{y})[-C(q, \beta)L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} + C(n, q, \beta)L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)-1} \\ & + C(n, q, \beta)L_2^2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}]. \end{aligned}$$

Now our aim is to make the term at the right-hand side of the above display to be strictly less than 0, through choosing L_1 large enough. We first select L_1 so large that

$$\begin{cases} \frac{1}{3}C(q, \beta)L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} \geq C(n, q, \beta)L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)-1}, \\ \frac{1}{3}C(q, \beta)L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} \geq C(n, q, \beta)L_2^2L_1^{q-3}|\bar{x} - \bar{y}|^{(\beta-1)(q-3)}, \end{cases}$$

i.e.,

$$L_1|\bar{x} - \bar{y}|^{\beta-1} \geq C(n, q, \beta, L_2).$$

This can be realized if L_1 is sufficiently large, since $|\bar{x} - \bar{y}| < 1$ and the power of it is negative. Next, we proceed to choose L_1 so large that

$$\begin{cases} \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq C(n, p, \beta)L_2L_1^{p-2}|\bar{x} - \bar{y}|^{(\beta-1)(p-2)-1}, \\ \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq C(n, p, \beta)L_2^2L_1^{p-3}|\bar{x} - \bar{y}|^{(\beta-1)(p-3)}, \\ \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq C(p, q, \beta, \gamma_\infty)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)}, \\ \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq C(n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty)L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)}, \\ \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq C(n, q, \beta, \|a\|_{C^1(B_1)})L_2L_1^{q-2}|\bar{x} - \bar{y}|^{(\beta-1)(q-2)+1}, \\ \frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} \geq \|\Phi\|_{L^\infty(B_1)}. \end{cases}$$

We arrange the previous display as

$$\begin{cases} L_1|\bar{x} - \bar{y}|^{\beta-1} \geq C(n, p, \beta, L_2, \|\Phi\|_{L^\infty(B_1)}), \\ |\bar{x} - \bar{y}|^{-1} \geq C(p, q, \beta, \gamma_\infty), \\ L_1^{p-q}|\bar{x} - \bar{y}|^{(\beta-1)(p-q)-1} \geq C(n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty), \\ L_1^{p-q+1}|\bar{x} - \bar{y}|^{(\beta-1)(p-q+1)-2} \geq C(n, p, q, \beta, \|a\|_{C^1(B_1)}, L_2). \end{cases} \tag{5.7}$$

Making use of the fact $|\bar{x} - \bar{y}| \leq \left(\frac{2\|u\|_{L^\infty(B_1)}}{L_1}\right)^{\frac{1}{\beta}}$, we can pick such large $L_1 > 1$ that the first two inequalities of (5.7) hold true. To assure the inequality (5.7)₃ holds, we first require

$$(\beta - 1)(p - q) - 1 < 0 \Rightarrow q < p + \frac{1}{1 - \beta} \quad (\beta \in (0, 1)).$$

We in turn enforce

$$\begin{aligned} L_1^{p-q}|\bar{x} - \bar{y}|^{(\beta-1)(p-q)-1} &\geq L_1^{p-q} \left(\frac{2\|u\|_{L^\infty(B_1)}}{L_1}\right)^{\frac{(\beta-1)(p-q)-1}{\beta}} \\ &\geq C(n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty), \end{aligned}$$

that is,

$$L_1^{\frac{p-q+1}{\beta}} \geq C(n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty, \|u\|_{L^\infty(B_1)}),$$

which is true precisely if

$$p - q + 1 > 0 \Rightarrow q < p + 1.$$

Under this condition, the inequality (5.7)₄ shall hold true by choosing L_1 large.

As has been shown above, if $q < p + 1$, we can select L_1 large enough, which depends on $n, p, q, \beta, \|a\|_{C^1(B_1)}, \gamma_\infty, \|u\|_{L^\infty(B_1)}$ and $\|\Phi\|_{L^\infty(B_1)}$, to get

$$0 \leq -\frac{1}{7}C(p, \beta)L_1^{p-1}|\bar{x} - \bar{y}|^{(\beta-1)(p-1)-1} - \frac{1}{3}C(q, \beta)a(\bar{y})L_1^{q-1}|\bar{x} - \bar{y}|^{(\beta-1)(q-1)-1} < 0.$$

This is a contradiction. Thus the claim (5.1) holds true, which means that the viscosity solution u is locally β -Hölder continuous. The proof is finished now. □

Based on the local β -Hölder continuity of u in Lemma 5.1, we could further deduce u is locally Lipschitz continuous.

Lemma 5.2 (Local Lipschitz continuity) *Let u be a bounded viscosity solution to (1.1) in B_1 . Under the assumptions that $0 \leq a(x) \in C^1(B_1)$, $p \leq q \leq p + \frac{1}{2}$ and (1.3), there is a constant C that depends on $n, p, q, \gamma_\infty, \|a\|_{C^1(B_1)}, \|u\|_{L^\infty(B_1)}$ and $\|\Phi\|_{L^\infty(B_1)}$, such that*

$$|u(x) - u(y)| \leq C|x - y|$$

for all $x, y \in B_{1/2}$. Here $\gamma_\infty := \max_{t \in [0, \|u\|_{L^\infty(B_1)}]} \gamma(t)$.

Proof Let $x_0, y_0 \in B_{1/2}$. Construct an auxiliary function

$$\begin{aligned} \Psi(x, y) &:= u(x) - u(y) - M_1\varphi(|x - y|) \\ &\quad - \frac{M_2}{2}|x - x_0|^2 - \frac{M_2}{2}|y - y_0|^2, \quad M_1, M_2 > 0, \end{aligned}$$

where

$$\varphi(t) := \begin{cases} t - \kappa_0 t^\nu, & \text{if } 0 \leq t \leq t_1 := \left(\frac{1}{4\nu\kappa_0}\right)^{\frac{1}{\nu-1}}, \\ \varphi(t_1), & \text{if } t > t_1 \end{cases}$$

with $1 < \nu < 2$ and $0 < \kappa_0 < 1$ such that $2 < t_1$. We are ready to verify $\Psi(x, y) \leq 0$ for $(x, y) \in B_{3/4} \times B_{3/4}$ under the appropriate choice of M_1, M_2 , which leads to Lipschitz continuity of u . We argue by contradiction. Suppose that Ψ reaches its positive maximum at $(\hat{x}, \hat{y}) \in \overline{B}_{3/4} \times \overline{B}_{3/4}$. As in the beginning of the proof of

Lemma 5.1, we can see that $\hat{x} \neq \hat{y}$ and $\hat{x}, \hat{y} \in B_{3/4}$ for large $M_2 \geq C\|u\|_{L^\infty(B_1)}$. Moreover, by Lemma 5.1, we know that

$$|u(x) - u(y)| \leq c_\beta |x - y|^\beta \quad \text{for } x, y \in B_{3/4},$$

where c_β is the same as the C in Lemma 5.1. From the assumptions, we further get

$$M_2|\hat{x} - x_0|, M_2|\hat{y} - y_0| \leq c_\beta |\hat{x} - \hat{y}|^{\frac{\beta}{2}} \tag{5.8}$$

by adjusting the constants (by letting $2M_2 \leq c_\beta$). Additionally, it follows, by selecting κ_0 sufficiently small, that

$$M_1(|\hat{x} - \hat{y}| - \kappa_0|\hat{x} - \hat{y}|^\nu) \leq 2\|u\|_{L^\infty(B_1)},$$

i.e.,

$$|\hat{x} - \hat{y}| \leq \frac{4\|u\|_{L^\infty(B_1)}}{M_1}. \tag{5.9}$$

Indeed, due to $|\hat{x} - \hat{y}| \leq 2$ and $\nu - 1 > 0$, we can fix $\kappa_0 \in (0, 1)$ such that $\frac{1}{2} \leq 1 - \kappa_0|\hat{x} - \hat{y}|^{\nu-1}$.

From the maximum principle for semicontinuous functions, for any $\mu > 0$, there exist $X, Y \in \mathcal{S}^n$ such that

$$(\eta_1, X + M_2I) \in \bar{J}^{2,+}u(\hat{x}) \quad \text{and} \quad (\eta_2, Y - M_2I) \in \bar{J}^{2,-}u(\hat{y})$$

and

$$-(\mu + 2\|B\|) \begin{pmatrix} I & \\ & I \end{pmatrix} \leq \begin{pmatrix} X & \\ & -Y \end{pmatrix} \leq \begin{pmatrix} B & -B \\ -B & B \end{pmatrix} + \frac{2}{\mu} \begin{pmatrix} B^2 & -B^2 \\ -B^2 & B^2 \end{pmatrix}, \tag{5.10}$$

where

$$\begin{aligned} \eta_1 &= M_1\varphi'(|\hat{x} - \hat{y}|) \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} + M_2(\hat{x} - x_0), \\ \eta_2 &= M_1\varphi'(|\hat{x} - \hat{y}|) \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} - M_2(\hat{y} - y_0) \end{aligned}$$

and

$$B = M_1\varphi''(|\hat{x} - \hat{y}|) \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} \otimes \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} + \frac{M_1\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \left(I - \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} \otimes \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|} \right).$$

Notice that, for $t \in [0, t_1]$,

$$\begin{cases} \varphi'(t) = 1 - \nu\kappa_0 t^{\nu-1}, \\ \varphi''(t) = -\nu(\nu - 1)\kappa_0 t^{\nu-2}, \end{cases} \tag{5.11}$$

and then $\frac{3}{4} \leq \varphi'(t) \leq 1$ and $\varphi''(t) < 0$ when $t \in (0, 2]$. Through straightforward calculation we get

$$\frac{M_1}{2} \leq |\eta_1|, |\eta_2| \leq 2M_1, \quad \text{if } M_1 \geq 4c_\beta, \tag{5.12}$$

$$\|B\| \leq M_1 \frac{\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \tag{5.13}$$

and

$$\|B^2\| \leq M_1^2 \left(|\varphi''(|\hat{x} - \hat{y}|)| + \frac{\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \right)^2. \tag{5.14}$$

According to (5.10), we obtain $X \leq Y$ and furthermore $\|X\|, \|Y\| \leq 2\|B\| + \mu$. Via taking

$$\mu = 4M_1 \left(|\varphi''(|\hat{x} - \hat{y}|)| + \frac{\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \right),$$

we have, for $\xi = \frac{\hat{x} - \hat{y}}{|\hat{x} - \hat{y}|}$,

$$\langle (X - Y)\xi, \xi \rangle \leq 4 \left(\langle B\xi, \xi \rangle + \frac{2}{\mu} \langle B^2\xi, \xi \rangle \right) \leq 2M_1\varphi''(|\hat{x} - \hat{y}|) < 0. \tag{5.15}$$

It follows from the last inequality that at lowest one eigenvalue of $X - Y$, denoted by $\hat{\lambda}$, is smaller than $2M_1\varphi''(|\hat{x} - \hat{y}|) < 0$. Besides, putting together (5.13), (5.14) and (5.10), we derive

$$\begin{aligned} \|Y\| &\leq 2|\langle B\bar{\xi}, \bar{\xi} \rangle| + \frac{4}{\mu} |\langle B^2\bar{\xi}, \bar{\xi} \rangle| \\ &\leq 4M_1 \left(|\varphi''(|\hat{x} - \hat{y}|)| + \frac{\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \right), \end{aligned} \tag{5.16}$$

where $\bar{\xi}$ is a unit vector.

Because u is a viscosity solution, we arrive at

$$\begin{aligned} 0 &\leq |\eta_1|^{p-2} \text{tr}(A_p(\eta_1)(X + M_2I)) - |\eta_2|^{p-2} \text{tr}(A_p(\eta_2)(Y - M_2I)) \\ &\quad + a(\hat{x})|\eta_1|^{q-2} \text{tr}(A_q(\eta_1)(X + M_2I)) - a(\hat{y})|\eta_2|^{q-2} \text{tr}(A_q(\eta_2)(Y - M_2I)) \\ &\quad + |\eta_1|^{q-2} \eta_1 \cdot Da(\hat{x}) - |\eta_2|^{q-2} \eta_2 \cdot Da(\hat{y}) + f(\hat{x}, u(\hat{x}), \eta_1) - f(\hat{y}, u(\hat{y}), \eta_2) \\ &=: T_1 + T_2 + T_3 + T_4. \end{aligned} \tag{5.17}$$

In the following, we want to prove $T_1 + T_2 + T_3 + T_4 < 0$ by taking M_1 large enough, the procedure of which is analogous to that in proof of Lemma 5.1. Here we shall briefly write down it. We first consider the term T_2 ,

$$\begin{aligned}
 T_2 &= (a(\hat{x}) - a(\hat{y}))|\eta_2|^{q-2}\text{tr}(A_q(\eta_2)(Y - M_2I)) + a(\hat{x})|\eta_1|^{q-2}\text{tr}(A_q(\eta_1)(X - Y)) \\
 &\quad + a(\hat{x})(|\eta_1|^{q-2} - |\eta_2|^{q-2})\text{tr}(A_q(\eta_2)Y) + a(\hat{x})|\eta_1|^{q-2}\text{tr}((A_q(\eta_1) - A_q(\eta_2))Y) \\
 &\quad + M_2a(\hat{x})[|\eta_1|^{q-2}\text{tr}(A_q(\eta_1)) + |\eta_2|^{q-2}\text{tr}(A_q(\eta_2))] \\
 &=: T_{21} + T_{22} + T_{23} + T_{24} + T_{25}.
 \end{aligned}$$

In view of (5.11), (5.12) and (5.16), there holds that

$$\begin{aligned}
 T_{21} &\leq C(q, \|a\|_{C^1(B_1)})M_1^{q-2}|\hat{x} - \hat{y}|(n\|A_q(\eta_2)\|\|Y\| + nqM_2) \\
 &\leq C(n, q, \|a\|_{C^1(B_1)})M_1^{q-2}[M_1(1 + |\hat{x} - \hat{y}|\varphi''(|\hat{x} - \hat{y}|)) + M_2] \\
 &\leq C(n, q, \|a\|_{C^1(B_1)})M_1^{q-2}[M_1(1 + |\hat{x} - \hat{y}|^{v-1}) + M_2] \\
 &\leq C(n, q, \|a\|_{C^1(B_1)})(M_1^{q-1} + M_1^{q-2}M_2).
 \end{aligned}$$

Applying the mean value theorem along with (5.8), (5.11) and (5.12), we obtain

$$\left| |\eta_1|^{q-2} - |\eta_2|^{q-2} \right| \leq C(q, c_\beta)M_1^{q-3}|\hat{x} - \hat{y}|^{\frac{\beta}{2}},$$

which indicates that

$$\begin{aligned}
 T_{23} &\leq C(n, q, c_\beta)a(\hat{x})M_1^{q-3}|\hat{x} - \hat{y}|^{\frac{\beta}{2}}\|Y\| \\
 &\leq C(n, c_\beta, q)a(\hat{x})M_1^{q-2}(|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + |\hat{x} - \hat{y}|^{v-2}).
 \end{aligned}$$

Thanks to (5.12) and (5.15), we derive

$$\begin{aligned}
 T_{22} &\leq a(\hat{x})|\eta_1|^{q-2} \sum_{i=1}^n \lambda_i(A_q(\eta_1))\lambda_i(X - Y) \\
 &\leq C(q)a(\hat{x})M_1^{q-1}\varphi''(|\hat{x} - \hat{y}|) \\
 &= -C(q, v, \kappa_0)a(\hat{x})M_1^{q-1}|\hat{x} - \hat{y}|^{v-2}.
 \end{aligned}$$

According to (5.8), (5.12) and (5.16),

$$\begin{aligned}
 T_{24} &\leq a(\hat{x})|\eta_1|^{q-2}n\|A_q(\eta_1) - A_q(\eta_2)\|\|Y\| \\
 &\leq C(n, q)a(\hat{x})|\eta_1|^{q-2} \max \left\{ \frac{|\eta_1 - \eta_2|}{|\eta_1|}, \frac{|\eta_1 - \eta_2|}{|\eta_2|} \right\} \|Y\| \\
 &\leq C(n, q, c_\beta)a(\hat{x})M_1^{q-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}} \left(|\varphi''(|\hat{x} - \hat{y}|)| + \frac{\varphi'(|\hat{x} - \hat{y}|)}{|\hat{x} - \hat{y}|} \right) \\
 &\leq C(n, q, c_\beta)a(\hat{x})M_1^{q-2}(|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + |\hat{x} - \hat{y}|^{v-2}).
 \end{aligned}$$

Finally, by (5.12) we have

$$T_{25} \leq C(n, q)a(\hat{x})M_2M_1^{q-2}.$$

Thus taking $\nu = \frac{\beta}{2} + 1$ and combining these previous inequalities yields that

$$T_2 \leq -C(q, \beta)a(\hat{x})M_1^{q-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, q, c_\beta)a(\hat{x})M_1^{q-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, q, \|a\|_{C^1(B_1)})(M_1^{q-1} + M_1^{q-2}M_2).$$

Similarly, we have

$$T_1 \leq -C(p, \beta)M_1^{p-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, p, c_\beta)M_1^{p-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, p)M_1^{p-2}M_2.$$

For T_3 , it is easy to get

$$T_3 \leq C(q, \|a\|_{C^1(B_1)})M_1^{q-1}.$$

Owing to the growth condition on f , we can see that

$$\begin{aligned} T_4 &\leq \gamma_\infty(|\eta_1|^{p-1} + a(\hat{x})|\eta_1|^{q-1}) + \gamma_\infty(|\eta_2|^{p-1} + a(\hat{y})|\eta_2|^{q-1}) + 2\|\Phi\|_{L^\infty(B_1)} \\ &\leq C(p, q, \gamma_\infty, \|a\|_{L^\infty(B_1)})(M_1^{p-1} + M_1^{q-1}) + 2\|\Phi\|_{L^\infty(B_1)} \\ &\leq C(p, q, \gamma_\infty, \|a\|_{L^\infty(B_1)}, \|\Phi\|_{L^\infty(B_1)})M_1^{q-1} \end{aligned}$$

with $\gamma_\infty := \max_{t \in [0, \|u\|_{L^\infty(B_1)}]} \gamma(t)$. Here we note $M_1 > 1$ is a sufficiently large number. It follows from merging the estimates on T_1 – T_4 with (5.17) that

$$\begin{aligned} 0 \leq & \left[-C(p, \beta)M_1^{p-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, p, c_\beta)M_1^{p-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \right. \\ & \left. + C(p, q, \gamma_\infty, \|a\|_{C^1(B_1)}, \|\Phi\|_{L^\infty(B_1)})M_1^{q-1} \right] \\ & + a(\hat{x}) \left[-C(q, \beta)M_1^{q-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} + C(n, q, c_\beta)M_1^{q-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \right], \end{aligned} \tag{5.18}$$

where we have used the relation $M_1 \geq M_2$ to simplify the display. To get a contradiction, we have to select such large M_1 that

$$\begin{cases} \frac{1}{3}C(p, \beta)M_1^{p-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \geq C(n, p, c_\beta)M_1^{p-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1}, \\ \frac{1}{3}C(p, \beta)M_1^{p-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \geq C(p, q, \gamma_\infty, \|a\|_{C^1(B_1)}, \|\Phi\|_{L^\infty(B_1)})M_1^{q-1}, \\ \frac{1}{2}C(q, \beta)M_1^{q-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \geq C(n, q, c_\beta)M_1^{q-2}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1}, \end{cases} \tag{5.19}$$

that is,

$$\begin{cases} M_1 \geq C(n, p, q, c_\beta), \\ M_1^{p-q}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \geq C(p, q, \beta, \gamma_\infty, \|a\|_{C^1(B_1)}, \|\Phi\|_{L^\infty(B_1)}). \end{cases}$$

Remembering (5.9), we arrive at

$$M_1^{p-q}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} \geq (4\|u\|_{L^\infty(B_1)})^{\frac{\beta}{2}-1}M_1^{p-q+1-\frac{\beta}{2}}.$$

Now enforcing

$$q < p + 1 - \frac{\beta}{2},$$

we could choose such large M_1 that

$$M_1 \geq C(n, p, q, \beta, \gamma_\infty, \|a\|_{C^1(B_1)}, \|\Phi\|_{L^\infty(B_1)}, \|u\|_{L^\infty(B_1)}),$$

which ensures (5.19) holds true. Then the display (5.18) becomes

$$0 \leq -\frac{1}{3}CM_1^{p-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} - \frac{1}{2}a(\hat{x})CM_1^{q-1}|\hat{x} - \hat{y}|^{\frac{\beta}{2}-1} < 0,$$

which is a contradiction. Let us mention that we fix β to be a specific number so that M_1 does not depend on β . Up to now, we have justified the local Lipschitz continuity of u . \square

Once Lemma 5.2 is proved in B_1 , then the case of a bounded domain Ω follows by covering arguments. Therefore, we finish the proof of Theorem 1.2.

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Data Availability Data sharing is not applicable to this article as no datasets were generated or analysed during the current study.

Declarations

Conflict of interest The authors declare that there is no conflict of interest. We also declare that this manuscript has no associated data.

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References

1. Ahmida, Y., Chlebicka, I., Gwiazda, P., Youssfi, A.: Gossez's approximation theorems in Musielak–Orlicz–Sobolev spaces. *J. Funct. Anal.* **275**(9), 2538–2571 (2018)
2. Attouchi, A., Ruosteenoja, E.: Remarks on regularity for p -Laplacian type equations in non-divergence form. *J. Differ. Equ.* **265**, 1922–1961 (2018)

3. Attouchi, A., Parviainen, M., Ruosteenoja, E.: $C^{1,\alpha}$ regularity for the normalized p -Poisson problem. *J. Math. Pures Appl.* **108**(4), 553–591 (2017)
4. Baasandorj, S., Byun, S.S., Oh, J.: Calderón–Zygmund estimates for generalized double phase problems. *J. Funct. Anal.* **279**(7), 108670, 57 pp (2020)
5. Balci, A.K., Diening, L., Surnachev, M.: New examples on Lavrentiev gap using fractals. *Calc. Var. Part. Differ. Equ.* **59**(5), Paper No. 180, 34 pp (2020)
6. Baroni, P., Colombo, M., Mingione, G.: Harnack inequalities for double phase functionals. *Nonlinear Anal.* **121**, 206–222 (2015)
7. Baroni, P., Colombo, M., Mingione, G.: Regularity for general functionals with double phase. *Calc. Var. Part. Differ. Equ.* **57**(2), Paper No. 62, 48 pp (2018)
8. Barrios, B., Medina, M.: Equivalence of weak and viscosity solutions in fractional nonhomogeneous problems. *Math. Ann.* **381**, 1979–2012 (2021)
9. Byun, S.S., Oh, J.: Global gradient estimates for non-uniformly elliptic equations. *Calc. Var. Part. Differ. Equ.* **56**(2), Paper No. 46, 36 pp (2017)
10. Chlebicka, I., De Filippis, C.: Removable sets in non-uniformly elliptic problems. *Ann. Mat. Pura Appl.* **199**, 619–649 (2020)
11. Chlebicka, I., Gwiazda, P., Świerczewska-Gwiazda, A., Wróblewska-Kamińska, A.: *Partial Differential Equations in Anisotropic Musielak–Orlicz Spaces*. Springer Monographs in Mathematics, Springer, Cham (2021)
12. Colasuonno, F., Squassina, M.: Eigenvalues for double phase variational integrals. *Ann. Mat. Pura Appl.* **195**, 1917–1959 (2016)
13. Colombo, M., Mingione, G.: Regularity for double phase variational problems. *Arch. Ration. Mech. Anal.* **215**, 443–496 (2015)
14. Colombo, M., Mingione, G.: Calderón–Zygmund estimates and non-uniformly elliptic operators. *J. Funct. Anal.* **270**, 1416–1478 (2016)
15. Crandall, M. G.: Viscosity solutions: a primer. viscosity solutions and applications (Montecatini Terme, 1995). In: *Lecture Notes in Math*, vol. 1660, Springer, Berlin, pp. 1–43 (1997)
16. Crandall, M.G., Ishii, H., Lions, P.-L.: User’s guide to viscosity solutions of second order partial differential equations. *Bull. Am. Math. Soc.* **27**(1), 1–67 (1992)
17. De Filippis, C., Mingione, G.: Nonuniformly elliptic Schauder theory. [arXiv:2201.07369](https://arxiv.org/abs/2201.07369)
18. De Filippis, C., Mingione, G.: A borderline case of Calderón–Zygmund estimates for non-uniformly elliptic problems. *St. Petersburg Math. J.* **31**(3), 82–115 (2019)
19. De Filippis, C., Mingione, G.: Lipschitz bounds and nonautonomous integrals. *Arch. Ration. Mech. Anal.* **242**(2), 973–1057 (2021)
20. De Filippis, C., Palatucci, G.: Hölder regularity for nonlocal double phase equations. *J. Differ. Equ.* **267**(1), 547–586 (2019)
21. Esposito, L., Leonetti, F., Mingione, G.: Sharp regularity for functionals with (p, q) growth. *J. Differ. Equ.* **204**(1), 5–55 (2004)
22. Fang, Y., Zhang, C.: On weak and viscosity solutions of nonlocal double phase equations. *Int. Math. Res. Not. IMRN*. <https://doi.org/10.1093/imrn/rnab351>
23. Fang, Y., Zhang, C.: Equivalence between distributional and viscosity solutions for the double-phase equation. *Adv. Calc. Var.* **15**(4), 811–829 (2022)
24. Fang, Y., Zhang, C.: Regularity for quasi-linear parabolic equations with nonhomogeneous degeneracy or singularity. *Calc. Var. Partial Differ. Equ.* **62**(1), Paper No. 2, 46 pp (2023)
25. Fang, Y., Rădulescu, V.D., Zhang, C.: Regularity of solutions to degenerate fully nonlinear elliptic equations with variable exponent. *Bull. Lond. Math. Soc.* **53**, 1863–1878 (2021)
26. Fang, Y., Rădulescu, V.D., Zhang, C., Zhang, X.: Gradient estimates for multi-phase problems in Campanato spaces. *Indiana Univ. Math. J.* **71**(3), 1079–1099 (2022)
27. Fonseca, I., Malý, J., Mingione, G.: Scalar minimizers with fractal singular sets. *Arch. Ration. Mech. Anal.* **172**, 295–307 (2004)
28. Harjulehto, P., Hästö, P.: *Orlicz Spaces and Generalized Orlicz Spaces*. Lecture Notes in Mathematics, vol. 2236. Springer, Cham (2019)
29. Hästö, P., Ok, J.: Maximal regularity for local minimizers of non-autonomous functionals. *J. Eur. Math. Soc. (JEMS)* **24**, 1285–1334 (2022)
30. Ishii, H.: On the equivalence of two notions of weak solutions, viscosity solutions and distribution solutions. *Funkcial. Ekvac.* **38**, 101–120 (1995)

31. Ishii, H., Lions, P.-L.: Viscosity solutions of fully nonlinear second-order elliptic partial differential equations. *J. Differ. Equ.* **83**(1), 26–78 (1990)
32. Julin, V., Juutinen, P.: A new proof for the equivalence of weak and viscosity solutions for the p -Laplace equation. *Comm. Part. Differ. Equ.* **37**(5), 934–946 (2012)
33. Juutinen, P., Lindqvist, P., Manfredi, J.: On the equivalence of viscosity solutions and weak solutions for a quasi-linear equation. *SIAM J. Math. Anal.* **33**(3), 699–717 (2001)
34. Juutinen, P., Lukkari, T., Parviainen, M.: Equivalence of viscosity and weak solutions for the $p(x)$ -Laplacian. *Ann. Inst. H. Poincaré Anal. Non Linéaire* **27**(6), 1471–1487 (2010)
35. Karppinen, A., Lee, M.: Hölder continuity of the minimizer of an obstacle problem with generalized Orlicz growth. *Int. Math. Res. Not. IMRN* **19**, 15313–15354 (2022)
36. Korvenpää, J., Kuusi, T., Lindgren, E.: Equivalence of solutions to fractional p -Laplace type equations. *J. Math. Pure. Appl.* **132**, 1–26 (2019)
37. Le, P.: Liouville results for double phase problems in \mathbb{R}^n . *Qual. Theory Dyn. Syst.* **21**, 59, 18 pp (2022)
38. Lindqvist, P.: Notes on the stationary p -Laplace equation. *SpringerBriefs in Mathematics*, Springer, Cham (2019)
39. Lions, P.-L.: Optimal control of diffusion processes and Hamilton-Jacobi-Bellman equations, Part 2: viscosity solutions and uniqueness. *Comm. Partial Differ. Equ.* **8**, 1229–1276 (1983)
40. Marcellini, P.: Regularity and existence of solutions of elliptic equations with p, q -growth conditions. *J. Differ. Equ.* **90**, 1–30 (1991)
41. Marcellini, P.: Growth conditions and regularity for weak solutions to nonlinear elliptic pdes. *J. Math. Anal. Appl.* **501**(1), Paper No. 124408, 32 pp (2021)
42. Medina, M., Ochoa, P.: On viscosity and weak solutions for non-homogeneous p -Laplace equations. *Adv. Nonlinear Anal.* **8**, 468–481 (2019)
43. Medina, M., Ochoa, P.: Equivalence of solutions for non-homogeneous $p(x)$ -Laplace equations. *Math. Eng.* **5**(2), 1–19 (2022)
44. Mingione, G., Rădulescu, V.D.: Recent developments in problems with nonstandard growth and nonuniform ellipticity. *J. Math. Anal. Appl.* **501**(1), Paper No. 125197, 41 pp (2021)
45. Musielak, J.: Orlicz Spaces and Modular Spaces. *Lecture Notes in Mathematics*, vol. 1034. Springer-Verlag, Berlin (1983)
46. Papageorgiou, N.S., Pudelko, A., Rădulescu, V.D.: Non-autonomous (p, q) -equations with unbalanced growth. *Math. Ann.* (2022). <https://doi.org/10.1007/s00208-022-02381-0>
47. Papageorgiou, N.S., Rădulescu, V.D., Repovš, D.: Existence and multiplicity of solutions for double-phase Robin problems. *Bull. Lond. Math. Soc.* **52**(3), 546–560 (2020)
48. Parviainen, M., Vázquez, J.-L.: Equivalence between radial solutions of different parabolic gradient-diffusion equations and applications. *Ann. Scuola Norm. Sup. Pisa Cl. Sci.* **21**, 303–359 (2020)
49. Siltakoski, J.: Equivalence of viscosity and weak solutions for the normalized $p(x)$ -Laplacian. *Calc. Var. Partial Differential Equations* **57**(4), Paper No. 95, 20 pp (2018)
50. Siltakoski, J.: Equivalence of viscosity and weak solutions for a p -parabolic equation. *J. Evol. Equ.* **21**, 2047–2080 (2021)
51. Zhikov, V.V.: Averaging of functionals of the calculus of variations and elasticity theory. *Izv. Akad. Nauk SSSR Ser. Mat.* **50**(4), 675–710 (1986)
52. Zhikov, V.V.: Lavrentiev phenomenon and homogenization for some variational problems. *C. R. Acad. Sci. Paris, Ser. I Math.* **316**(5), 435–439 (1993)