

Uniqueness of the blow-up boundary solution of logistic equations with absorbtion *

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Abstract. Let Ω be a smooth bounded domain in \mathbf{R}^N . Assume $f \in C^1[0, \infty)$ is a non-negative function such that $f(u)/u$ is increasing on $(0, \infty)$. Let a be a real number and let $b \geq 0$, $b \not\equiv 0$ be a continuous function such that $b \equiv 0$ on $\partial\Omega$. We study the logistic equation $\Delta u + au = b(x)f(u)$ in Ω . The special feature of this work is the uniqueness of positive solutions blowing-up on $\partial\Omega$, in a general setting that arises in probability theory.

Unicité de la solution explosant au bord pour équations logistiques avec absorption

Résumé. Soit Ω un domaine borné et régulier de \mathbf{R}^N . On suppose que $f \in C^1[0, \infty)$ est une fonction non-négative telle que $f(u)/u$ soit strictement croissante sur $(0, +\infty)$. Soit a un réel et $b \geq 0$, $b \not\equiv 0$, une fonction continue sur $\bar{\Omega}$ telle que $b \equiv 0$ sur $\partial\Omega$. On étudie l'équation logistique $\Delta u + au = b(x)f(u)$ sur Ω . Le but de cette Note est de montrer l'unicité de la solution explosant au bord de Ω dans un contexte général, qui apparaît en théorie des probabilités.

Version française abrégée. Soit $\Omega \subset \mathbf{R}^N$ ($N \geq 3$) un domaine borné et régulier, a un paramètre réel et $b \in C^{0,\mu}(\bar{\Omega})$, $\mu \in (0, 1)$, $b \geq 0$, $b \not\equiv 0$ dans Ω . On considère l'équation logistique

$$\Delta u + au = b(x)f(u) \quad \text{dans } \Omega, \quad (1)$$

où $f \in C^1[0, \infty)$ satisfait

(A₁) $f \geq 0$ et $f(u)/u$ est strictement croissante sur $(0, +\infty)$.

Soit

$$\Omega_0 := \text{int} \{x \in \Omega : b(x) = 0\}$$

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et on suppose que $\partial\Omega_0$ est régulier (éventuellement vide), $\overline{\Omega}_0 \subset \Omega$ et $b > 0$ sur $\Omega \setminus \overline{\Omega}_0$. On désigne par $\lambda_{\infty,1}$ la première valeur propre (avec conditions de Dirichlet) de l'opérateur $(-\Delta)$ dans Ω_0 , avec la convention $\lambda_{\infty,1} = +\infty$ si $\Omega_0 = \emptyset$.

On dit que u est une solution *large (explosive)* de (1) si $u \geq 0$ dans Ω et $u(x) \rightarrow \infty$ si $d(x) := \text{dist}(x, \partial\Omega) \rightarrow 0$.

Soit $D > 0$ et $R : [D, \infty) \rightarrow (0, +\infty)$ une fonction mesurable. On dit que R a une variation régulière d'indice $\rho \in \mathbf{R}$ (notation: $R \in \mathbf{R}_\rho$) si $\lim_{u \rightarrow \infty} R(\xi u)/R(u) = \xi^\rho$, pour chaque $\xi > 0$ (voir [11]).

Soit \mathcal{K} l'ensemble des fonctions $k : (0, \nu) \rightarrow (0, +\infty)$ (pour un certain ν), de classe C^1 , croissantes, telles que $\lim_{t \rightarrow 0+} \left(\frac{\int_0^t k(s) ds}{k(t)} \right)^{(i)} := \ell_i$, pour $i = \overline{0,1}$.

On démontre le résultat suivant.

THÉORÈME 1. - *Supposons que la fonction f satisfait la condition (A_1) et que f' est une fonction à variation régulière d'indice $\rho \neq 0$. De plus, on suppose que le potentiel b vérifie*

(B) $b(x) = c k^2(d(x)) + o(k^2(d(x)))$ si $d(x) \rightarrow 0$, avec $c > 0$ et $k \in \mathcal{K}$.

Alors, pour chaque $a \in (-\infty, \lambda_{\infty,1})$, l'équation (1) admet une unique solution explosive u_a . On a, de plus,

$$\lim_{d(x) \rightarrow 0} \frac{u_a(x)}{h(d(x))} = \xi_0,$$

où $\xi_0 = \left(\frac{2 + \ell_1 \rho}{c(2 + \rho)} \right)^{1/\rho}$ et la fonction h est définie par

$$\int_{h(t)}^{\infty} \frac{ds}{\sqrt{2F(s)}} = \int_0^t k(s) ds, \quad \forall t \in (0, \nu).$$

Let $\Omega \subset \mathbf{R}^N$ ($N \geq 3$) be a smooth bounded domain. Consider the semilinear elliptic equation

$$\Delta u + au = b(x)f(u) \quad \text{in } \Omega, \tag{1}$$

where a is a real parameter and $b \in C^{0,\mu}(\overline{\Omega})$, for some $\mu \in (0, 1)$, such that $b \geq 0$, $b \not\equiv 0$ in Ω .

Suppose that $f \in C^1[0, \infty)$ satisfies

(A_1) $f \geq 0$ and $f(u)/u$ is increasing on $(0, \infty)$.

In the study of positive solutions for (1), subject to the homogeneous Dirichlet boundary condition, an important role is played by the zero set (see [1])

$$\Omega_0 := \text{int} \{x \in \Omega : b(x) = 0\}.$$

We shall assume throughout that Ω_0 is smooth (possibly empty), $\overline{\Omega}_0 \subset \Omega$, and $b > 0$ in $\Omega \setminus \overline{\Omega}_0$.

By a *large (explosive)* solution of (1) we mean a solution u of (1) such that $u \geq 0$ in Ω and $u(x) \rightarrow \infty$ as $d(x) := \text{dist}(x, \partial\Omega) \rightarrow 0$. In [3, 4] we study the existence of large solutions for (1)

and also deduce several existence and unicity results for a related problem. Note that any large solution of (1) is *positive* and it can exist only if the Keller-Osserman condition holds (see [4])

$$(A_2) \quad \int_1^\infty \frac{dt}{\sqrt{F(t)}} < \infty, \quad \text{where } F(t) = \int_0^t f(s) ds.$$

Let H_∞ define the Dirichlet Laplacian on the set $\Omega_0 \subset \Omega$ as the unique self-adjoint operator associated to the quadratic form $\psi(u) = \int_{\Omega} |\nabla u|^2 dx$ with form domain

$$H_D^1(\Omega_0) = \{u \in H_0^1(\Omega) : u(x) = 0 \text{ for a.e. } x \in \Omega \setminus \Omega_0\}.$$

If $\partial\Omega_0$ satisfies an exterior cone condition, then $H_D^1(\Omega_0)$ coincides with $H_0^1(\Omega_0)$ and H_∞ is the classical Laplace operator with Dirichlet condition on $\partial\Omega_0$.

Let $\lambda_{\infty,1}$ be the first Dirichlet eigenvalue of H_∞ in Ω_0 . We understand $\lambda_{\infty,1} = +\infty$ if $\Omega_0 = \emptyset$.

The main result in [3] asserts that equation (1) has a large solution iff $a \in (-\infty, \lambda_{\infty,1})$.

The special feature of this paper is the uniqueness of large solutions of (1) in a general framework for f and b , under the restriction $b \equiv 0$ on $\partial\Omega$, inherited from the logistic equation (see [6]).

We start with

DEFINITION 1 ([11]). - A positive measurable function R defined on $[D, \infty)$, for some $D > 0$, is called *regularly varying (at infinity) with index* $q \in \mathbf{R}$, written $R \in \mathbf{R}_q$, if for all $\xi > 0$

$$\lim_{u \rightarrow \infty} R(\xi u)/R(u) = \xi^q.$$

When the index of regular variation q is zero, we say that the function is *slowly varying*.

REMARK 1. - Any function $R \in \mathbf{R}_q$ can be written in terms of a slowly varying function. Indeed, set $R(u) = u^q L(u)$. From Definition 1 we easily derive that L varies slowly.

The canonical q -varying function is u^q . The functions $\ln(1+u)$, $\ln \ln(e+u)$, $\exp\{(\ln u)^\alpha\}$, $\alpha \in (0, 1)$ vary slowly, as well as any measurable function on $[D, \infty)$ with positive limit at infinity.

In what follows L denotes an arbitrary slowly varying function and $D > 0$ a positive number. For details on Properties 1-4 stated below, we refer to Seneta [11] (pp. 7, 18, 53 and 78).

PROPERTY 1. - For any $m > 0$, $u^m L(u) \rightarrow \infty$, $u^{-m} L(u) \rightarrow 0$ as $u \rightarrow \infty$.

PROPERTY 2. - Any positive C^1 -function on $[D, \infty)$ satisfying $uL_1'(u)/L_1(u) \rightarrow 0$ as $u \rightarrow \infty$ is slowly varying. Moreover, if the above limit is $q \in \mathbf{R}$, then $L_1 \in \mathbf{R}_q$.

PROPERTY 3. - Assume $R : [D, \infty) \rightarrow (0, \infty)$ is measurable and Lebesgue integrable on each finite subinterval of $[D, \infty)$. Then R varies regularly iff there exists $j \in \mathbf{R}$ such that

$$\lim_{u \rightarrow \infty} \frac{u^{j+1} R(u)}{\int_D^u x^j R(x) dx} \quad (2)$$

exists and is a positive number, say $a_j + 1$. In this case, $R \in \mathbf{R}_q$ with $q = a_j - j$.

PROPERTY 4 (Karamata Theorem, 1933). - If $R \in \mathbf{R}_q$ is Lebesgue integrable on each finite subinterval of $[D, \infty)$, then the limit defined by (2) is $q + j + 1$, for every $j > -q - 1$.

LEMMA 1. - Assume (A_1) holds. Then we have the equivalence

$$a) f' \in \mathbf{R}_p \iff b) \lim_{u \rightarrow \infty} u f'(u)/f(u) := \vartheta < \infty \iff c) \lim_{u \rightarrow \infty} (F/f)'(u) := \gamma > 0.$$

REMARK 2. - Let $a)$ of Lemma 1 be fulfilled. The following assertions hold

- (i) ρ is **non-negative**. Indeed, if $\rho < 0$ then Property 1 and Remark 1 would contradict (A_1) ;
- (ii) $\gamma = 1/(\rho + 2) = 1/(\vartheta + 1)$ (see the proof of Lemma 1);
- (iii) If $\rho \neq 0$, then (A_2) holds (use $\lim_{u \rightarrow \infty} f(u)/u^p = \infty, \forall p \in (1, 1 + \rho)$). The converse implication is not necessarily true (take $f(u) = u \ln^4(u + 1)$). However, there are cases when $\rho = 0$ and (A_2) fails so that (1) has **no** large solutions. This is illustrated by $f(u) = u$ or $f(u) = u \ln(u + 1)$.

Inspired by the definition of γ , we denote by \mathcal{K} the set of all positive, increasing C^1 -functions k defined on $(0, \nu)$, for some $\nu > 0$, which satisfy $\lim_{t \rightarrow 0+} \left(\frac{\int_0^t k(s) ds}{k(t)} \right)^{(i)} := \ell_i, i = \overline{0, 1}$.

It is easy to see that $\ell_0 = 0$ and $\ell_1 \in [0, 1]$, for every $k \in \mathcal{K}$. Our next result gives examples of functions $k \in \mathcal{K}$ with $\lim_{t \rightarrow 0+} k(t) = 0$, for every $\ell_1 \in [0, 1]$.

LEMMA 2. - Let $S \in C^1[D, \infty)$ be such that $S' \in \mathbf{R}_q$ with $q > -1$. Hence the following hold:

- a) If $k(t) = \exp \{-S(1/t)\} \quad \forall t \leq 1/D$, then $k \in \mathcal{K}$ with $\ell_1 = 0$.
- b) If $k(t) = 1/S(1/t) \quad \forall t \leq 1/D$, then $k \in \mathcal{K}$ with $\ell_1 = 1/(q + 2) \in (0, 1)$.
- c) If $k(t) = 1/\ln S(1/t) \quad \forall t \leq 1/D$, then $k \in \mathcal{K}$ with $\ell_1 = 1$.

REMARK 3. - If $S \in C^1[D, \infty)$, then $S' \in \mathbf{R}_q$ with $q > -1$ iff for some $m > 0, C > 0$ and $B > D$ we have $S(u) = Cu^m \exp \left\{ \int_B^u \frac{y(t)}{t} dt \right\}, \forall u \geq B$, where $y \in C[B, \infty)$ satisfies $\lim_{u \rightarrow \infty} y(u) = 0$. In this case, $S' \in \mathbf{R}_q$ with $q = m - 1$. This is a consequence of Properties 3 and 4.

Our main result is

THEOREM 1. - Let (A_1) hold and $f' \in \mathbf{R}_\rho$ with $\rho > 0$. Assume $b \equiv 0$ on $\partial\Omega$ satisfies

- (B) $b(x) = c k^2(d(x)) + o(k^2(d(x)))$ as $d(x) \rightarrow 0$, for some constant $c > 0$ and $k \in \mathcal{K}$.

Then, for any $a \in (-\infty, \lambda_{\infty, 1})$, Eq. (1) admits a unique large solution u_a . Moreover,

$$\lim_{d(x) \rightarrow 0} \frac{u_a(x)}{h(d(x))} = \xi_0, \quad (3)$$

where $\xi_0 = \left(\frac{2 + \ell_1 \rho}{c(2 + \rho)} \right)^{1/\rho}$ and h is defined by

$$\int_{h(t)}^{\infty} \frac{ds}{\sqrt{2F(s)}} = \int_0^t k(s) ds, \quad \forall t \in (0, \nu). \quad (4)$$

By Remark 3, the assumption $f' \in \mathbf{R}_\rho$ with $\rho > 0$ holds iff there exist $p > 1$ and $B > 0$ such that $f(u) = Cu^p \exp \left\{ \int_B^u \frac{y(t)}{t} dt \right\}$, for all $u \geq B$ (y as before and $p = \rho + 1$). If B is large enough ($y > -\rho$ on $[B, \infty)$), then $f(u)/u$ is increasing on $[B, \infty)$. Thus, to get the whole range of functions f for which our Theorem 1 applies we have only to “paste” a suitable smooth function on $[0, B]$ in accordance with (A_1) . A simple way to do this is to define $f(u) = u^p \exp \left\{ \int_0^u \frac{z(t)}{t} dt \right\}$, for all $u \geq 0$, where $z \in C[0, \infty)$ is non-negative such that $\lim_{t \rightarrow 0+} z(t)/t \in [0, \infty)$ and $\lim_{u \rightarrow \infty} z(u) = 0$. Clearly, $f(u) = u^p$, $f(u) = u^p \ln(u + 1)$, and $f(u) = u^p \arctan u$ ($p > 1$) fall into this category.

Lemma 2 provides a practical method to find functions k which can be considered in the statement of Theorem 1. Here are some examples: $k(t) = \exp \{-1/t^\alpha\}$, $k(t) = \exp \{-\ln(1 + \frac{1}{t})/t^\alpha\}$, $k(t) = \exp \{-[\arctan(\frac{1}{t})]/t^\alpha\}$, $k(t) = -1/\ln t$, $k(t) = t^\alpha/\ln(1 + \frac{1}{t})$, $k(t) = t^\alpha$, for some $\alpha > 0$.

As we shall see, the uniqueness lies upon the crucial observation (3), which shows that all explosive solutions have the same boundary behaviour. Note that the only case of Theorem 1 studied so far is $f(u) = u^p$ ($p > 1$) and $k(t) = t^\alpha$ ($\alpha > 0$) (see [6]). For related results on the uniqueness of explosive solutions (mainly in the cases $b \equiv 1$ and $a = 0$) we refer to [2, 8, 9, 12].

Proof of Lemma 1. - From Property 4 and Remark 2 (i) we deduce $a) \implies b)$ and $\vartheta = \rho + 1$. Conversely, $b) \implies a)$ follows by Property 3 since $\vartheta \geq 1$ cf. (A_1) .

$b) \implies c)$. Indeed, $\lim_{u \rightarrow \infty} \frac{uf(u)}{F(u)} = 1 + \vartheta$, which yields $\frac{\vartheta}{1+\vartheta} = \lim_{u \rightarrow \infty} \left[1 - \left(\frac{F}{f} \right)'(u) \right] = 1 - \gamma$.

$c) \implies b)$. Choose $s_1 > 0$ such that $\left(\frac{F}{f} \right)'(u) \geq \frac{\gamma}{2}$, $\forall u \geq s_1$. So, $\left(\frac{F}{f} \right)'(u) \geq \frac{(u-s_1)\gamma}{2} + \left(\frac{F}{f} \right)'(s_1)$, $\forall u \geq s_1$. Passing to the limit $u \rightarrow \infty$, we find $\lim_{u \rightarrow \infty} \frac{F(u)}{f(u)} = \infty$. Thus, $\lim_{u \rightarrow \infty} \frac{uf(u)}{F(u)} = \frac{1}{\gamma}$. Since $1 - \gamma := \lim_{u \rightarrow \infty} \frac{F(u)f'(u)}{f^2(u)}$, we obtain $\lim_{u \rightarrow \infty} \frac{uf'(u)}{f(u)} = \frac{1-\gamma}{\gamma}$. ■

Proof of Lemma 2. - Since $\lim_{u \rightarrow \infty} uS'(u) = \infty$ (cf. Property 1), from Karamata Theorem we deduce $\lim_{u \rightarrow \infty} \frac{uS'(u)}{S(u)} = q + 1 > 0$. Therefore, in any of the cases $a)$, $b)$, $c)$, $\lim_{t \rightarrow 0+} k(t) = 0$ and k is an increasing C^1 -function on $(0, \nu)$, for $\nu > 0$ sufficiently small.

$a)$ It is clear that $\lim_{t \rightarrow 0+} \frac{tk'(t)}{k(t) \ln k(t)} = \lim_{t \rightarrow 0+} \frac{-S'(1/t)}{tS(1/t)} = -(q+1)$. By l'Hospital's rule, $\ell_0 = \lim_{t \rightarrow 0+} \frac{k(t)}{k'(t)} = 0$ and $\lim_{t \rightarrow 0+} \frac{\left(\int_0^t k(s) ds \right) \ln k(t)}{tk(t)} = -\frac{1}{q+1}$. So, $1 - \ell_1 := \lim_{t \rightarrow 0+} \frac{\left(\int_0^t k(s) ds \right) k'(t)}{k^2(t)} = 1$.

$b)$ We see that $\lim_{t \rightarrow 0+} \frac{tk'(t)}{k(t)} = \lim_{t \rightarrow 0+} \frac{S'(1/t)}{tS(1/t)} = q + 1$. By l'Hospital's rule, $\ell_0 = 0$ and $\lim_{t \rightarrow 0+} \frac{\int_0^t k(s) ds}{tk(t)} = \frac{1}{q+2}$. So, $\ell_1 = 1 - \lim_{t \rightarrow 0+} \frac{\int_0^t k(s) ds}{tk(t)} \frac{tk'(t)}{k(t)} = \frac{1}{q+2}$.

$c)$ We have $\lim_{t \rightarrow 0+} \frac{tk'(t)}{k^2(t)} = \lim_{t \rightarrow 0+} \frac{S'(1/t)}{tS(1/t)} = q + 1$. By l'Hospital's rule, $\lim_{t \rightarrow 0+} \frac{\int_0^t k(s) ds}{tk(t)} = 1$. Thus, $\ell_0 = 0$ and $\ell_1 = 1 - \lim_{t \rightarrow 0+} \frac{\int_0^t k(s) ds}{t} \frac{tk'(t)}{k^2(t)} = 1$. ■

Proof of Theorem 1. - Fix $a \in (-\infty, \lambda_{\infty,1})$. By [3, Theorem 1], (1) has at least a large solution.

If we prove that (3) holds for an *arbitrary* large solution u_a of (1), then the uniqueness is a consequence of [3, Lemma 3]. Indeed, if u_1 and u_2 are two arbitrary large solutions of (1), then (3) yields $\lim_{d(x) \rightarrow 0+} \frac{u_1(x)}{u_2(x)} = 1$. Hence, for any $\varepsilon \in (0, 1)$, there exists $\delta = \delta(\varepsilon) > 0$ such that

$$(1 - \varepsilon)u_2(x) \leq u_1(x) \leq (1 + \varepsilon)u_2(x), \quad \forall x \in \Omega \text{ with } 0 < d(x) \leq \delta. \quad (5)$$

Choosing eventually a smaller $\delta > 0$, we can assume that $\overline{\Omega}_0 \subset C_\delta$, where $C_\delta := \{x \in \Omega : d(x) > \delta\}$.

It is clear that u_1 is a positive solution of the boundary value problem

$$\Delta \phi + a\phi = b(x)f(\phi) \quad \text{in } C_\delta, \quad \phi = u_1 \quad \text{on } \partial C_\delta. \quad (6)$$

By (A_1) and (5), we see that $\phi^- = (1 - \varepsilon)u_2$ (resp., $\phi^+ = (1 + \varepsilon)u_2$) is a positive sub-solution (resp., super-solution) of (6). By the sub and super-solutions method, (6) has a positive solution ϕ_1 satisfying $\phi^- \leq \phi_1 \leq \phi^+$ in C_δ . Since $b > 0$ on $\overline{C_\delta} \setminus \overline{\Omega}_0$, by [3, Lemma 3] we derive that (6) has a *unique* positive solution, i.e., $u_1 \equiv \phi_1$ in C_δ . This yields $(1 - \varepsilon)u_2(x) \leq u_1(x) \leq (1 + \varepsilon)u_2(x)$ in C_δ , so that (5) holds in Ω . Passing to the limit $\varepsilon \rightarrow 0^+$, we conclude that $u_1 \equiv u_2$.

In order to prove (3) we state some useful properties about h :

(h_1) $h \in C^2(0, \nu)$, $\lim_{t \rightarrow 0+} h(t) = \infty$ (straightforward from (4)).

(h_2) $\lim_{t \rightarrow 0+} \frac{h''(t)}{k^2(t)f(h(t)\xi)} = \frac{1}{\xi^{\rho+1}} \frac{2 + \rho\ell_1}{2 + \rho}$, $\forall \xi > 0$ (so, $h'' > 0$ on $(0, 2\delta)$, for $\delta > 0$ small enough).

(h_3) $\lim_{t \rightarrow 0+} h(t)/h''(t) = \lim_{t \rightarrow 0+} h'(t)/h''(t) = 0$.

We check (h_2) for $\xi = 1$ only, since $f \in \mathbf{R}_{\rho+1}$. Clearly, $h'(t) = -k(t)\sqrt{2F(h(t))}$ and

$$h''(t) = k^2(t)f(h(t)) \left(1 - 2 \frac{k'(t) \left(\int_0^t k(s) ds \right)}{k^2(t)} \frac{\sqrt{F(h(t))}}{f(h(t)) \int_{h(t)}^\infty [F(s)]^{-1/2} ds} \right) \quad \forall t \in (0, \nu). \quad (7)$$

We see that $\lim_{u \rightarrow \infty} \sqrt{F(u)}/f(u) = 0$. Thus, from l'Hospital's rule and Lemma 1 we infer that

$$\lim_{u \rightarrow \infty} \frac{\sqrt{F(u)}}{f(u) \int_u^\infty [F(s)]^{-1/2} ds} = \frac{1}{2} - \gamma = \frac{\rho}{2(\rho + 2)}. \quad (8)$$

Using (7) and (8) we derive (h_2) and also

$$\lim_{t \rightarrow 0^+} \frac{h'(t)}{h''(t)} = \frac{-2(2 + \rho)}{2 + \ell_1 \rho} \lim_{t \rightarrow 0^+} \frac{\int_0^t k(s) ds}{k(t)} \lim_{u \rightarrow \infty} \frac{\sqrt{F(u)}}{f(u) \int_u^\infty [F(s)]^{-1/2} ds} = \frac{-\rho \ell_0}{2 + \ell_1 \rho} = 0. \quad (9)$$

From (h_1) and (h_2) , $\lim_{t \rightarrow 0^+} h'(t) = -\infty$. So, l'Hospital's rule and (9) yield $\lim_{t \rightarrow 0^+} \frac{h(t)}{h'(t)} = 0$. This and (9) lead to $\lim_{t \rightarrow 0^+} \frac{h(t)}{h''(t)} = 0$ which proves (h_3) .

Proof of (3). Fix $\varepsilon \in (0, c/2)$. Since $b \equiv 0$ on $\partial\Omega$ and (B) holds, we take $\delta > 0$ so that

- (i) $d(x)$ is a C^2 -function on the set $\{x \in \mathbf{R}^N : d(x) < 2\delta\}$;
- (ii) k^2 is increasing on $(0, 2\delta)$;
- (iii) $(c - \varepsilon)k^2(d(x)) < b(x) < (c + \varepsilon)k^2(d(x))$, $\forall x \in \Omega$ with $0 < d(x) < 2\delta$;
- (iv) $h''(t) > 0 \forall t \in (0, 2\delta)$ (from (h_2)).

Let $\sigma \in (0, \delta)$ be arbitrary. We define $\xi^\pm = \left[\frac{2 + \ell_1 \rho}{(c \mp 2\varepsilon)(2 + \rho)} \right]^{1/\rho}$ and $v_\sigma^-(x) = h(d(x) + \sigma)\xi^-$, for all x with $d(x) + \sigma < 2\delta$ resp., $v_\sigma^+(x) = h(d(x) - \sigma)\xi^+$, for all x with $\sigma < d(x) < 2\delta$.

Using (i)-(iv), when $\sigma < d(x) < 2\delta$ we obtain (since $|\nabla d(x)| \equiv 1$)

$$\Delta v_\sigma^+ + a v_\sigma^+ - b(x)f(v_\sigma^+) \leq \xi^+ h''(d(x) - \sigma) \left(\frac{h'(d(x) - \sigma)}{h''(d(x) - \sigma)} \Delta d(x) + a \frac{h(d(x) - \sigma)}{h''(d(x) - \sigma)} + 1 - (c - \varepsilon) \frac{k^2(d(x) - \sigma)f(h(d(x) - \sigma)\xi^+)}{h''(d(x) - \sigma)\xi^+} \right).$$

Similarly, when $d(x) + \sigma < 2\delta$ we find

$$\Delta v_\sigma^- + a v_\sigma^- - b(x)f(v_\sigma^-) \geq \xi^- h''(d(x) + \sigma) \left(\frac{h'(d(x) + \sigma)}{h''(d(x) + \sigma)} \Delta d(x) + a \frac{h(d(x) + \sigma)}{h''(d(x) + \sigma)} + 1 - (c + \varepsilon) \frac{k^2(d(x) + \sigma)f(h(d(x) + \sigma)\xi^-)}{h''(d(x) + \sigma)\xi^-} \right).$$

Using (h_2) and (h_3) we see that, by diminishing δ , we can assume

$$\Delta v_\sigma^+(x) + a v_\sigma^+(x) - b(x)f(v_\sigma^+(x)) \leq 0 \quad \forall x \text{ with } \sigma < d(x) < 2\delta;$$

$$\Delta v_\sigma^-(x) + a v_\sigma^-(x) - b(x)f(v_\sigma^-(x)) \geq 0 \quad \forall x \text{ with } d(x) + \sigma < 2\delta.$$

Let Ω_1 and Ω_2 be smooth bounded domains such that $\Omega \subset \subset \Omega_1 \subset \subset \Omega_2$ and the first Dirichlet eigenvalue of $(-\Delta)$ in the domain $\Omega_1 \setminus \overline{\Omega}$ is greater than a . Let $p \in C^{0,\mu}(\overline{\Omega}_2)$ satisfy $0 < p(x) \leq b(x)$ for $x \in \Omega \setminus C_{2\delta}$, $p = 0$ on $\overline{\Omega}_1 \setminus \Omega$ and $p > 0$ on $\Omega_2 \setminus \overline{\Omega}_1$. Denote by w a positive large solution of

$$\Delta w + a w = p(x)f(w) \quad \text{in } \Omega_2 \setminus \overline{C}_{2\delta}.$$

The existence of w is ensured by Theorem 1 in [3].

Suppose that u_a is an arbitrary large solution of (1) and let $v := u_a + w$. Then v satisfies

$$\Delta v + av - b(x)f(v) \leq 0 \quad \text{in } \Omega \setminus \overline{C}_{2\delta}.$$

Since $v|_{\partial\Omega} = \infty > v_{\sigma}^-|_{\partial\Omega}$ and $v|_{\partial C_{2\delta}} = \infty > v_{\sigma}^-|_{\partial C_{2\delta}}$, Lemma 1 in [3] implies

$$u_a + w \geq v_{\sigma}^- \quad \text{on } \Omega \setminus \overline{C}_{2\delta}. \quad (10)$$

Similarly,

$$v_{\sigma}^+ + w \geq u_a \quad \text{on } C_{\sigma} \setminus \overline{C}_{2\delta}. \quad (11)$$

Letting $\sigma \rightarrow 0$ in (10) and (11), we deduce $h(d(x))\xi^+ + 2w \geq u_a + w \geq h(d(x))\xi^-$, for all $x \in \Omega \setminus \overline{C}_{2\delta}$.

Since w is uniformly bounded on $\partial\Omega$, we have $\xi^- \leq \liminf_{d(x) \rightarrow 0} \frac{u_a(x)}{h(d(x))} \leq \limsup_{d(x) \rightarrow 0} \frac{u_a(x)}{h(d(x))} \leq \xi^+$.

Letting $\varepsilon \rightarrow 0^+$ we obtain (3). This concludes the proof of Theorem 1. \blacksquare

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